§10.5-10.6 Partial differential equations, separation of variables.

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Partial differential equations, distinctive features; simple examples.

Possibly the simplest PDE that we can imagine is

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The general solution is

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Partial differential equations, distinctive features; simple examples.

Possibly the simplest PDE that we can imagine is

$$\frac{\partial u(x,t)}{\partial t} = 0$$

The general solution is

$$u(x,t) = f(x)$$

for *any* function *f*!

Solutions of PDEs always have "functional degree of freedom" (as opposed to free constants in the case of ODEs).



$$\frac{\partial u(x,t)}{\partial t} = \frac{\partial u(x,t)}{\partial x} \tag{1}$$

Claim: the general solution is



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Claim: the general solution is f(x + t), for any differentiable function f. Indeed, by the chain rule

$$\frac{\partial f(x+t)}{\partial t} = f'(x+t) = \frac{\partial f(x+t)}{\partial x}$$



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Exercise: Show that there are no other solutions!



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Exercise: Show that there are no other solutions!What do we need to specify to determine the particular solution we are interested in? Since we have functional degree of freedom, we need to specify, say an initial **function**.

$\partial u(x,t)$	 $\partial u(x,t)$
∂t	 ∂x



$$\frac{\partial u(x,t)}{\partial t} = \frac{\partial u(x,t)}{\partial x}$$

For instance (2) can be solved uniquely if we specify
 $u(x,0) = U(x)$.

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(2)

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 $u(x,0) = U(x)$. Say, we give $u(x,0) = \sin^3(x)$. Then indeed, the
general solution $f(x + t)$ must satisfy the condition
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$$u_n(x, t) = c_n \exp(-n^2 \alpha^2 \pi^2 t/L^2) \sin(n\pi x/L)$$

is a particular solution for any *n*, and any constant c_n .



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$$\alpha^2 u_{xx} = c_n \alpha^2 (-(n\pi/L)^2 \sin(n\pi x/L)) \exp(-n^2 \alpha^2 \pi^2 t/L^2)$$



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$$\alpha^2 u_{xx} = c_n \alpha^2 (-(n\pi/L)^2 \sin(n\pi x/L)) \exp(-n^2 \alpha^2 \pi^2 t/L^2)$$

3. So:
$$u_t = \alpha^2 u_{xx}$$
 for all x and t .



$$u_n(x,t) = c_n \exp(-n^2 \alpha^2 \pi^2 t/L^2) \sin(n\pi x/L)$$

4. Boundary conditions? Check: u(0, t) = u(L, t) = 0.

5. Initial condition? $u(x, 0) = c_n \sin(n\pi x/L)$. This is not general enough. We need more solutions.





$$u_t = \alpha^2 u_{xx}$$

1

Find all solutions of the form

$$u(x,t) = X(x)T(t):$$



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 so

$$\frac{T'(t)}{\alpha^2 T(t)}$$
depends on *t* alone

$$\underbrace{\frac{X''(x)}{X(x)}}_{\text{depends on } x \text{ alone}}$$

$$\blacksquare \blacksquare \diamondsuit \models \models \longleftarrow \longleftarrow \rightarrow$$











 $\frac{T'(t)}{\alpha^2 T(t)} = \frac{X''(x)}{X(x)}$ depends on t alone depends on x alone
Thus $\frac{X''(x)}{X(x)}$ is simply a constant, say $-\lambda$. Then $\frac{T'(t)}{\alpha^2 T(t)}$ is equal to the same constant.

We arrive at a pair of ODEs:

$$\frac{T'(t)}{\alpha^2 T(t)} = -\lambda$$

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We arrive at a pair of ODEs:

$$\frac{T'(t)}{\alpha^2 T(t)} = -\lambda \tag{3}$$

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(5)

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The **boundary value problem** (6) is an eigenvalue problem

$$X''(x) = -\lambda X(x); \quad X(0) = 0, \ X(L) = 0$$

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The **boundary value problem** (6) is an eigenvalue problem

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(7)

where we seek **nonzero solutions**.


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(8)

We studied (8) before.



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, $n = 1, 2, 3, ...$



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and the eigenfunctions are

 $X_n = (c_n)\sin(n\pi x/L)$



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Putting X_n and T_n together –remember,

 $u_n(x, t) = X_n(x)T_n(t)$ we have:

$$u_n(x,t) = c_n \exp(-n^2 \alpha^2 \pi^2 t/L^2) \sin(n\pi x/L)$$



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(9)

is also a solution of the problem. (We'll deal with convergence: later.)





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Can this be now made to fit any initial temperature distribution, u(x, 0) = U(x)? Yes, by the Fourier theorem.



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That ensures that



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That ensures that (1) U_1 has a pure sine FS. (2) $U_1 = U$ on the interval of interest, [0, L].





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$$U_1(x) = \begin{cases} -20 \text{ for } x \in (-L, 0) \\ 20 \text{ for } x \in (0, L) \end{cases}$$



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$$\frac{1}{L} \int_0^L U_1 \sin(n\pi x/L) dx = \frac{1}{L} \int_0^L 20 \sin(n\pi x/L) dx = 40 \frac{1 - (-1)^n}{n\pi}$$

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The complete solution is thus

$$u(x,t) = \sum_{n=1}^{\infty} c_n \exp(-n^2 \pi^2 \alpha^2 t / L^2) \sin(n\pi x / L)$$

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Other Heat Equation settings

Nonhomogeneous boundary conditions. Here, we seek to solve

$$u_t = \alpha^2 u_{xx}, \quad u(0, t) = T_1, u(L, t) = T_2, u(x, 0) = f(x)$$



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that is, we have different temperatures at the endpoints. As in nonhomogeneous ODEs, the solution is essentially **any solution of the nonhomogeneous equation plus the general solution of the homogeneous one**.



Indeed if u_0 satisfies the eq, boundary conditions but not necessarily the initial condition, then if we write $u = u_0 + v$ we have $(u_0)_t + v_t = \alpha^2 (u_0)_{xx} + v_{xx}$ or $v_t + \underbrace{((u_0)_t - (u_0)_{xx})}_{=0,by \text{ construction}} = \alpha^2 v_{xx}$

We need $v(0, t) + u_0(0, t) = T_1$ but $u_0(0, t) = T_1$, by construction, so: v(0, t) = 0. Likewise, v(L, t) = 0. v satisfies the same problem, with homogneous boundary values, and initial condition

$$v(x,0) + u_0(x,0) = f(x) \implies$$

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$$v(x,0) + u_0(x,0) = f(x) \implies v(x,0) = f(x) - u_0(x,0)$$



A particular solution of

$$u_t = \alpha^2 u_{xx}, \quad u(0,t) = T_1, u(L,t) = T_2$$

is easy to find. Look, for instance for solutions that don't depend on *t*. Then

 $u_{xx}=0$,



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 $u_{xx} = 0, \Rightarrow u = Ax + B;$


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