## §10.5-10.6 Partial differential

 equations, separation of variables.
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O. Costin: §10.4-5

## Partial differential equations, distinctive features; simple

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The general solution is

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u(x, t)=f(x)
$$

for any function $f$ !
Solutions of PDEs always have "functional degree of freedom" (as opposed to free constants in the case of ODEs).

Another simple example:

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\frac{\partial u(x, t)}{\partial t}=\frac{\partial u(x, t)}{\partial x} \tag{1}
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Claim: the general solution is $f(x+t)$, for any differentiable function $f$. Indeed, by the chain rule

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Exercise: Show that there are no other solutions! What do we need to specify to determine the particular solution we are interested in? Since we have functional degree of freedom, we need to specify, say an initial function.

$$
\frac{\partial u(x, t)}{\partial t}=\frac{\partial u(x, t)}{\partial x}
$$

$$
\begin{equation*}
\frac{\partial u(x, t)}{\partial t}=\frac{\partial u(x, t)}{\partial x} \tag{2}
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For instance (2) can be solved uniquely if we specify $u(x, 0)=U(x)$.

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u_{n}(x, t)=c_{n} \exp \left(-n^{2} \alpha^{2} \pi^{2} t / L^{2}\right) \sin (n \pi x / L)
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is a particular solution for any $n$, and any constant $c_{n}$. What we mean by $u_{n}$ being a solution of the PDE is precisely that $\left(u_{n}(x, t)\right)_{t}=\alpha^{2}\left(u_{n}(x, t)\right)_{x x}$ for any $t$ and $x$.

1. $u_{t}=-\left(n^{2} \alpha^{2} \pi^{2} / L^{2}\right) \exp \left(-n^{2} \alpha^{2} \pi^{2} t / L^{2}\right) c_{n} \sin (n \pi x / L)$

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2. $\alpha^{2} u_{x x}=c_{n} \alpha^{2}\left(-(n \pi / L)^{2} \sin (n \pi x / L)\right) \exp \left(-n^{2} \alpha^{2} \pi^{2} t / L^{2}\right)$

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3. So: $u_{t}=\alpha^{2} u_{x x}$ for all $x$ and $t$.

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4. Boundary conditions? Check: $u(0, t)=u(L, t)=0$.
5. Initial condition? $u(x, 0)=c_{n} \sin (n \pi x / L)$. This is not general enough. We need more solutions.

Separation of variables.
O. Costin: §10.4-5

Separation of variables.

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Find all solutions of the form

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$$
X(x) T^{\prime}(t)=\alpha^{2} X^{\prime \prime}(x) T(t) \text { so } \underbrace{\frac{T^{\prime}(t)}{\alpha^{2} T(t)}}_{\text {depends on } t \text { alone }}=\underbrace{\frac{X^{\prime \prime}(x)}{X(x)}}_{\text {depends on } x \text { alone }}
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We arrive at a pair of ODEs:

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\frac{T^{\prime}(t)}{\alpha^{2} T(t)}=-\lambda
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\frac{T^{\prime}(t)}{\alpha^{2} T(t)} & =-\lambda  \tag{3}\\
\frac{X^{\prime \prime}(x)}{X(x)} & =-\lambda \tag{4}
\end{align*}
$$

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O. Costin: §10.4-5

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where we seek nonzero solutions.
O. Costin: §10.4-5

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O. Costin: §10.4-5

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O. Costin: §10.4-5
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Putting $X_{n}$ and $T_{n}$ together -remember,

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\begin{gathered}
u_{n}(x, t)=X_{n}(x) T_{n}(t) \quad \text { we have: } \\
u_{n}(x, t)=c_{n} \exp \left(-n^{2} \alpha^{2} \pi^{2} t / L^{2}\right) \sin (n \pi x / L)
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O. Costin: §10.4-5

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u(x, t)=\sum_{n=1}^{\infty} c_{n} \exp \left(-n^{2} \pi^{2} \alpha^{2} t / L^{2}\right) \sin (n \pi x / L)
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\end{equation*}
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is also a solution of the problem. (We'll deal with convergence: later.)

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## Initial condition We have

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\begin{equation*}
u(x, 0)=\sum_{n=1}^{\infty} c_{n} \sin (n \pi x / L) \tag{11}
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Can this be now made to fit any initial temperature distribution, $u(x, 0)=U(x)$ ? Yes, by the Fourier theorem.

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That ensures that (1) $U_{1}$ has a pure sine FS. (2) $U_{1}=U$ on the interval of interest, $[0, L]$.

$$
f(x)=\sum_{n=1}^{\infty} c_{n} \exp \left(-n^{2} \pi^{2} \alpha^{2} 0 / L^{2}\right) \sin (n \pi x / L)=\sum_{n=1}^{\infty} c_{n} \sin (n \pi x / L)
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In our example $U(x)=20$, thus the function to be worked with is

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U_{1}(x)=\left\{\begin{array}{l}
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\frac{1}{L} \int_{0}^{L} U_{1} \sin (n \pi x / L) d x=\frac{1}{L} \int_{0}^{L} 20 \sin (n \pi x / L) d x=40 \frac{1-(-1)^{n}}{n \pi}
$$

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f(x)=\sum_{n=1}^{\infty} c_{n} \exp \left(-n^{2} \pi^{2} \alpha^{2} 0 / L^{2}\right) \sin (n \pi x / L)=\sum_{n=1}^{\infty} c_{n} \sin (n \pi x / L)
$$

In our example $U(x)=20$, thus the function to be worked with is

$$
U_{1}(x)=\left\{\begin{array}{l}
-20 \text { for } x \in(-L, 0)  \tag{12}\\
20 \text { for } x \in(0, L)
\end{array}\right.
$$

Since this is indeed an odd function, the coefficients $c_{n}$ are given by

$$
\frac{1}{L} \int_{0}^{L} U_{1} \sin (n \pi x / L) d x=\frac{1}{L} \int_{0}^{L} 20 \sin (n \pi x / L) d x=40 \frac{1-(-1)^{n}}{n \pi}
$$

The complete solution is thus

$$
u(x, t)=\sum_{n=1}^{\infty} c_{n} \exp \left(-n^{2} \pi^{2} \alpha^{2} t / L^{2}\right) \sin (n \pi x / L)
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O. Costin: §10.4-5

## Other Heat Equation settings

Nonhomogeneous boundary conditions. Here, we seek to solve

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u_{t}=\alpha^{2} u_{x x}, \quad u(0, t)=T_{1}, u(L, t)=T_{2}, u(x, 0)=f(x)
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that is, we have different temperatures at the endpoints. As in nonhomogeneous ODEs, the solution is essentially any solution of the nonhomogeneous equation plus the general solution of the homogeneous one.

Indeed if $u_{0}$ satisfies the eq, boundary conditions but not necessarily the initial condition, then if we write $u=u_{0}+v$ we have $\left(u_{0}\right)_{t}+v_{t}=\alpha^{2}\left(u_{0}\right)_{x x}+v_{x x}$ or $v_{t}+\underbrace{\left(\left(u_{0}\right)_{t}-\left(u_{0}\right)_{x x}\right)}_{=0, \text { by construction }}=\alpha^{2} v_{x x}$

We need $v(0, t)+u_{0}(0, t)=T_{1}$ but $u_{0}(0, t)=T_{1}$, by construction, so: $v(0, t)=0$. Likewise, $v(L, t)=0 . v$ satisfies the same problem, with homogneous boundary values, and initial condition

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v(x, 0)+u_{0}(x, 0)=f(x) \Longrightarrow
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v(x, 0)+u_{0}(x, 0)=f(x) \Longrightarrow v(x, 0)=f(x)-u_{0}(x, 0)
$$

A particular solution of

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u_{t}=\alpha^{2} u_{x x}, \quad u(0, t)=T_{1}, u(L, t)=T_{2}
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is easy to find. Look, for instance for solutions that don't depend on $t$. Then

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& u_{x x}=0, \Rightarrow u=A x+B ; \quad A 0+B=T_{1}, A L+B=T_{2} \\
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$v_{t}=\alpha^{2} v_{x x}, \quad v(0, t)=0, v(L, t)=0, v(x, 0)=f(x)-\left[x\left(T_{2}-T_{1}\right) / L+T_{1}\right]$

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