

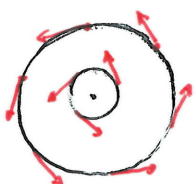
§1 Curl of a general rotational vector field:

Special vector fields: $\vec{F}_{(x,y,z)} = \vec{a} \times \vec{r}$ where $\vec{a} = \langle a_1, a_2, a_3 \rangle$ is a constant vector, $\vec{a} \neq \vec{0}$, & $\vec{r} = \langle x, y, z \rangle$. In coordinates:

$$\vec{F} = \vec{a} \times \vec{r} = \begin{vmatrix} i & j & k \\ a_1 & a_2 & a_3 \\ x & y & z \end{vmatrix} = \langle a_2z - a_3y, a_3x - a_1z, a_1y - a_2x \rangle$$

Def: \vec{F} as above is called a general rotational field in \mathbb{R}^3 .

Eg: $a_1 = a_2 = 0, a_3 = 1$ we get $\vec{F} = \langle -y, x, 0 \rangle = \langle -y, x \rangle$ usual rotational vector field in \mathbb{R}^2



concentric circles: $|\vec{F}|$ is constant & \vec{F} is tangent to the circles.

axis of rotation: positive z-axis. curl = $a_x - b_y = 1 - (-1) = 2 > 0$

& counterclockwise rotation looking from above.

In general: $\vec{F} = \vec{a} \times \vec{r}$ is a rotational v. field with axis of rotation in the direction of \vec{a}

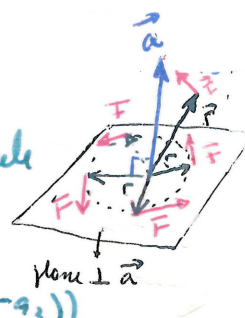
3 Properties ① $\nabla \cdot \vec{F} = \frac{\partial}{\partial x}(a_2z - a_3y) + 0 + 0 = 0$, so $\text{div } \vec{F} = 0$

② $\vec{F} = \vec{a} \times \vec{r}$ circles, the vector \vec{a} in the counterclockwise direction looking along \vec{a} from head to tail.

Why? \vec{F} lies in an orthogonal plane to \vec{a} & use right-hand rule

$$\textcircled{3} \text{curl}(\vec{F}) = \nabla \times \vec{F} = \begin{vmatrix} i & j & k \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ f & g & h \end{vmatrix} = i(a_1 - (-a_1)) - j(-a_2 - a_2) + k(a_3 - (-a_3)) = \langle 2a_1, 2a_2, 2a_3 \rangle = 2\vec{a}$$

$$\begin{cases} f = a_2z - a_3y \\ g = a_3x - a_1z \\ h = a_1y - a_2x \end{cases}$$

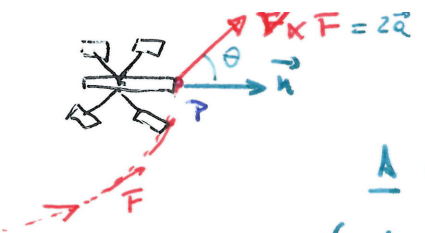


So $\text{curl}(\vec{F})$ has the same direction as the axis of rotation of \vec{F} (\vec{a}).

Furthermore $|\text{curl}(\vec{F})| = |\nabla \times \vec{F}| = 2|\vec{a}|$

Here, $|\vec{a}|$ is the constant angular speed of rotation = rate at which a particle rotates about the axis of rotation of the v. field \vec{F} .

Remark Suppose a paddle wheel is placed in \vec{F} at a pt P with the axis of the wheel in the direction of a unit vector \vec{n} . Q: For which \vec{n} 's does the paddle spin faster?



$$\text{comp}_{\vec{n}} \nabla \times \vec{F} = \frac{(\nabla \times \vec{F}) \cdot \vec{n}}{|\vec{n}|} = \frac{|\nabla \times \vec{F}| \cdot |\vec{n}| \cos \theta}{|\vec{n}|} = |\nabla \times \vec{F}| \cos \theta$$

Want to maximize this! So $\cos \theta = 1$, i.e. $\theta = 0$

Conclusion: Fastest rotation speed when \vec{n} has the same dir as $\nabla \times \vec{F}$.

If $\theta = \pm \frac{\pi}{2}$, then no rotation at all (the wheel doesn't spin).

§2 Interpreting the curl via Stokes' Thm: Write $C = \text{curve bounding the surface } S$

Flex integral along a surface = net flow of the vector field across the surface
 (eg: water flowing through a pipe's cross section)

If $\vec{G} = \langle f, g, h \rangle$ is a vector field & P is a pt in a surface S in \mathbb{R}^3 w/ unit outer normal \vec{n} perpendicular to the tangent plane to S at P .

Flex at P is the scalar component of $\vec{G}_{(P)}$ in the direction of \vec{n} .

$$\text{comp}_{\vec{n}} \vec{G}_{(P)} = \frac{\vec{G}_{(P)} \cdot \vec{n}}{|\vec{n}|} = \frac{|\vec{G}_{(P)}| |\vec{n}| \cos \theta}{|\vec{n}|} = |\vec{G}_{(P)}| \cos \theta$$

- If $\theta = 0 \Rightarrow$ all of $\vec{G}_{(P)}$ flows across S in the dir. of \vec{n} .
- If $\theta = \pi \Rightarrow$ opposite to \vec{n}
- If $\theta = \frac{\pi}{2} \Rightarrow \vec{G} \cdot \vec{n} = 0$ no flow through S at P .

\Rightarrow Flex integral adds up these scalar components. to compute the net flow of \vec{G} across S .

Assume the unit outer normal to S has the same dir as $t_u \times t_v = \frac{\partial \vec{r}}{\partial u} \times \frac{\partial \vec{r}}{\partial v}$.

Take any vector field \vec{F} , then $\vec{G} = \text{curl}(\vec{F}) = \nabla \times \vec{F}$ is a new vector field.

Stokes' Thm: Assume \vec{r} is 1-to-1 in the interior of R & with 1st partials in an open set in the uv -plane containing R . Assume R is cont & simply connected & bounded by a simple closed piecewise smooth curve C' . Then

$$\text{circ}(F, C) = \oint_C \vec{F} \cdot d\vec{r} = \iint_S (\nabla \times \vec{F}) \cdot \vec{n} \, dS = \text{Flex}(\nabla \times \vec{F}, S)$$

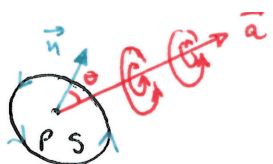
\vec{n} is unit outer normal

We can use this to define the average flux of $\nabla \times \vec{F}$ over S :

$$\text{Def Average flux of } \nabla \times \vec{F} \text{ over } S = \frac{1}{\text{area}(S)} \oint_C \vec{F} \cdot d\vec{r} = \frac{1}{\text{area}(S)} \iint_S (\nabla \times \vec{F}) \cdot \vec{n} \, dS$$

Special case: $\vec{F} = \vec{a} \times \vec{r}$ general rot. field in \mathbb{R}^3 .

Know: $\nabla \times \vec{F} = 2\vec{a}$.



Take S = small circular disk centered at a point P w/ normal \vec{n} forming an angle θ with \vec{a} .

Take C = boundary circle, counterclockwise oriented from above

$$\boxed{\text{Average flux}} = \frac{1}{\text{area}(S)} \iint_S (\underbrace{\nabla_x \vec{F}}_{\text{constant} = 2|\vec{a}| \cos \theta}) \cdot \vec{n} \, dS = \frac{1}{\text{area}(S)} (\nabla_x \vec{F}) \cdot \vec{n} \underbrace{\iint_S 1 \, dS}_{\text{area}(S)} = \boxed{2|\vec{a}| \cos \theta}$$

• If $\theta = 0 \Rightarrow$ average circ has max value $2|\vec{a}| = 2$ angular speed.

• If $\theta = \frac{\pi}{2} \Rightarrow$ " " is 0.

Conclusions: • The scalar comp of $\nabla_x \vec{F}$ at P in the dir of $\vec{n} =$ average flux of $\nabla_x \vec{F}$ on S .

• dir $(\nabla_x \vec{F})$ at P is the direction that maximizes the average flux of \vec{F} on S .

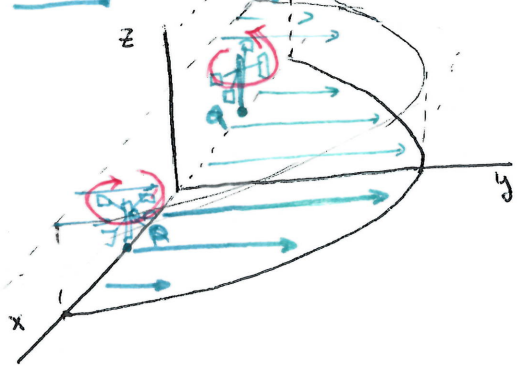
Equiv, we should orient the axis of a paddle in this direction to maximize the angular speed.

• General v. fields given P , take S_R small disk centered at a pt P w/ radius R & call $C_R =$ boundary of S_R

$$\lim_{R \rightarrow 0} (\text{Average flux}(\vec{F}, S_R)) = \lim_{R \rightarrow 0^+} \frac{1}{\text{area}(S_R)} \iint_{S_R} (\nabla_x \vec{F}) \cdot \vec{n} \, dS = (\nabla_x \vec{F}) \cdot \vec{n}_{(P)} \text{ fixed number.}$$

closed to $\nabla_x \vec{F}_{(P)} \cdot \vec{n}$ if R is small.

Example: Horizontal channel flow $\vec{v} = \langle 0, 1-x^2, 0 \rangle$ velocity v. field $|x|, |z| \leq 1$.



Place paddle wheels at $P = (\frac{1}{2}, 0, 0)$ & $Q = (-\frac{1}{2}, 0, 0)$.

Q: Find axis for each paddle wheel that makes it spin fastest.

- If in x -axis \Rightarrow NO spin! (flow strikes the upper & lower halves of the wheel symmetrically!).
- If in y -axis, w. spin! (flow parallel to axis!).
- If in z -axis: ① flow for $x < \frac{1}{2}$ greater than for $x > \frac{1}{2}$ so COUNTERCLOCKWISE SPIN at P .
- ② flow for $x < -\frac{1}{2}$ is smaller than for $x > -\frac{1}{2}$ so COUNTERCLOCKWISE SPIN at Q .

$$\text{Curl} = ? \quad \nabla_x \vec{v} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ 0 & 1-x^2 & 0 \end{vmatrix} = \langle 0, 0, -2x \rangle \text{ has } z\text{-axis direction.}$$

\Rightarrow max spin in this dir!

$$\vec{n} = \mathbf{k} = \langle 0, 0, 1 \rangle$$

$$(\nabla_x \vec{v}) \cdot \vec{n} = -2x > 0 \text{ at } Q \Rightarrow \text{counterclockwise spin!}$$

$$(\nabla_x \vec{v}) \cdot \vec{n} = -2x < 0 \text{ at } P \Rightarrow \text{clockwise spin!}$$

§ 3 Consequences of Stokes' Theorem: $\vec{F} = \langle f, g, h \rangle$

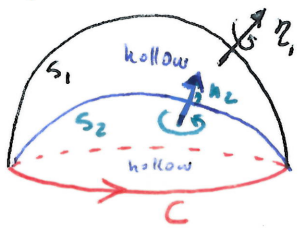
Corollary 1: If $f_y = g_x$, $f_z = h_x$ & $g_z = h_y$ in an open, connected & simply conn. region D of \mathbb{R}^3 , then $\oint_C \vec{F} \cdot d\vec{r} = 0$ on all closed simple sm. curves C in D & so \vec{F} is conservative \checkmark

Proof: Given C as above, we can find a smooth oriented surface S in D bounded by C . Note that $\nabla_x \vec{F} = \vec{0}$ by hypothesis, so Stoke's Theorem gives

$$\oint_C \vec{F} \cdot d\vec{r} = \iint_S (\nabla_x \vec{F}) \cdot \vec{n} \, dS = 0 \checkmark$$

Extensions: A curve C can bound more than 1 surface & we can use Stokes' Theorem to simplify the calculation of $\text{circ}(\vec{F}, C)$ or $\text{Flux}(\nabla_x \vec{F}, S_1)$ using $\text{Flux}(\nabla_x \vec{F}, S_2)$

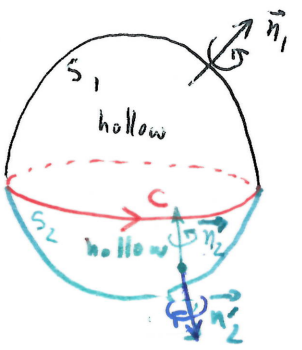
CASE 1:



\vec{n}_1 & \vec{n}_2 both point upwards.

$$\iint_{S_1} (\nabla_x \vec{F}) \cdot \vec{n}_1 \, dS = \oint_C \vec{F} \cdot d\vec{r} = \iint_{S_2} \nabla_x \vec{F} \cdot \vec{n}_2 \, dS$$

CASE 2:



\vec{n}_1 & \vec{n}_2' point in opposite directions & \vec{n}_2' is NOT comp. with the orientation of C . But \vec{n}_1 & \vec{n}_2' are comp. w/ orientation of S_2 .

$$\iint_{S_1} \nabla_x \vec{F} \cdot \vec{n}_1 \, dS = \oint_C \vec{F} \cdot d\vec{r}$$

$$\iint_{S_2} \nabla_x \vec{F} \cdot \vec{n}_2' \, dS = \oint_{C^{op}} \vec{F} \cdot d\vec{r} = -\oint_C \vec{F} \cdot d\vec{r} = -\iint_{S_2} \nabla_x \vec{F} \cdot \vec{n}_2 \, dS$$

S_1 & S_2 form a closed surface S w/ common boundary curve C

$$\oint_S \nabla_x \vec{F} \cdot \vec{n} \, dS = \iint_{S_2} \nabla_x \vec{F} \cdot \vec{n}_2' \, dS + \iint_{S_1} \nabla_x \vec{F} \cdot \vec{n}_1 \, dS = -\oint_C \vec{F} \cdot d\vec{r} + \oint_C \vec{F} \cdot d\vec{r} = \boxed{0}$$

Conclusion: On any closed, oriented, smooth surface S_1 , $\iint_{S_1} \nabla_x \vec{F} \cdot \vec{n} \, dS = 0$ enclosing a region in \mathbb{R}^3