

Recall: Laplace method for determining leading order behaviour

$$\text{of } I(x) = \int_a^b f(t) e^{-x\phi(t)} dt$$

$$\sim f(c) e^{-x\phi(c)} \sqrt{\frac{2\pi}{x\phi''(c)}} \quad \text{as } x \rightarrow \infty$$

if $c \in (a, b)$ is the unique min of ϕ on $[a, b]$, $\phi''(c) > 0$ and $f(c) \neq 0$.

§1. To determine full asymptotic expansion, we need Watson's

Lemma

$$\left[\begin{array}{l} f(t) \sim \sum_{n=0}^{\infty} c_n t^{a_n} \text{ as } t \rightarrow 0^+ \\ \Rightarrow \int_0^{\infty} f(t) e^{-xt} dt \sim \sum_{n=0}^{\infty} c_n \Gamma(a_n+1) x^{-a_n-1} \text{ as } x \rightarrow \infty. \end{array} \right]$$

Thus, in order to obtain higher order terms in Laplace's formula, we restrict ourselves to a neighbourhood of c (assume $c=0$ for simplicity).

$$\phi(t) - \phi(0) = \frac{\phi''(0)}{2} t^2 + \dots \text{ is 2-1 map near 0}$$

so, it has 2 local inverses: on $(-\epsilon, 0)$ and on $(0, \epsilon)$

$$I(x) \sim \underbrace{\int_{-\epsilon}^0 f(t) e^{-x\phi(t)} dt}_{I_-(x)} + \underbrace{\int_0^{\epsilon} f(t) e^{-x\phi(t)} dt}_{I_+(x)}$$

$$I_+(x) = \int_0^{\infty} f(t) e^{-x(\phi(t)-\phi(0))} \cdot e^{-x\phi(0)} dt$$

$w = \phi(t) - \phi(0)$ can be inverted

$$= \frac{\phi''(0)}{2} t^2 + \dots$$

$$t = a(w) + \sqrt{w} b(w) \\ = \sum_{n=1}^{\infty} a_n w^n + \sum_{m=0}^{\infty} b_m w^{m+\frac{1}{2}}$$

say, $t = h(w)$

$$dt = h'(w) dw$$

$$\left(b_0^2 = \frac{2}{\phi''(0)} \right)$$

$$\text{So } I_+(x) = \int_0^{\varepsilon' = \phi(\varepsilon) - \phi(0)} \underbrace{f(h(w)) \cdot h'(w)} e^{-xw} dw$$

$$\text{as } w \rightarrow 0, \quad \left(f(0) + w^{1/2}(\dots) \right) \left(\frac{b_0}{2} w^{-1/2} + \dots \right)$$

Watson's lemma applies. The problem becomes that of inverting a power series.

$$I_+(x) \sim f(0) \frac{1}{\sqrt{2\phi''(0)}} \Gamma\left(\frac{1}{2}\right) x^{-\frac{1}{2}} \quad \text{as } x \rightarrow \infty.$$

Repeating the same argument for $I_-(x) = \int_{-\varepsilon}^0 f(t) e^{-x\phi(t)} dt$

$$\text{gives } I(x) \sim f(0) \sqrt{\frac{2}{x\phi''(c)}} \cdot \Gamma\left(\frac{1}{2}\right).$$

Thus, we can obtain Laplace's formula from Watson's lemma.

§2. Stirling Series again.

$$\text{Recall } \Gamma(x+1) = \int_0^{\infty} t^x e^{-t} dt$$

$$= x^{x+1} \int_0^{\infty} s^x e^{-xs} ds \quad \left(\begin{array}{l} \text{set } t = x \cdot s \\ dt = x \cdot ds \end{array} \right)$$

$$\frac{\Gamma(x+1)}{x^{x+1}} = \int_0^{\infty} e^{-x(s - \ln(s))} ds$$

$$\phi(s) = s - \ln(s) \text{ has unique min. at } s=1$$

$$\left(\begin{array}{l} \phi'(s) = 1 - \frac{1}{s} < 0 \text{ for } s \in (0,1) \\ > 0 \text{ for } s \in (1,\infty) \end{array} \right) \quad \left(\begin{array}{l} \phi(1) = 1 \\ \phi''(1) = 1 \end{array} \right)$$

$$I(x) = \int_0^{\infty} e^{-x(s - \ln(s))} ds \sim \underbrace{\int_{1-\epsilon}^1 e^{-x(s - \ln(s))} ds}_{-I_-(x)} + \underbrace{\int_1^{1+\epsilon} e^{-x(s - \ln(s))} ds}_{I_+(x)}$$

Focus on $I_+(x)$:

$$I_+(x) = \int_0^{\epsilon} e^{-x(t+1 - \ln(t+1))} dt$$

$$= e^{-x} \int_0^{\epsilon} e^{-x(t - \ln(t+1))} dt$$

$$w = t - \ln(1+t) = \frac{t^2}{2} - \frac{t^3}{3} + \frac{t^4}{4} - \dots$$

can be (locally) inverted to $t = h(w) = \sum_{n=1}^{\infty} a_n w^n + \sum_{m=0}^{\infty} b_m w^{m+\frac{1}{2}}$ ($b_0 = \sqrt{2}$)

$$dt = h'(w) dw = \left(\sum_{n=1}^{\infty} n a_n w^{n-1} + \sum_{m=0}^{\infty} (m+\frac{1}{2}) b_m w^{m-\frac{1}{2}} \right) dw$$

$$I_+(x) = e^{-x} \int_0^{\varepsilon'} h'(w) e^{-xw} dw$$

$$\sim e^{-x} \left(\sum_{m=0}^{\infty} (m+\frac{1}{2}) b_m \Gamma(m+\frac{1}{2}) x^{-m-\frac{1}{2}} + \sum_{n=1}^{\infty} n a_n (n-1)! x^{-n} \right)$$

Same analysis as above gives $I_-(x)$ from $t = \bar{h}(w) = \sum a_n w^n - \sum b_m w^{m+\frac{1}{2}}$ (check)

$$\text{So, } I(x) \sim 2 e^{-x} \left(\sum_{m=0}^{\infty} b_m \Gamma(m+\frac{3}{2}) x^{-m-1} \right)$$

Exercise: Solve for first few terms $b_0, a_1, b_1, a_2, b_2, \dots$

Hint: $h'(w) = \frac{dt}{dw} = \frac{1+h(w)}{h(w)}$; write $h(w) = a(w) + \sqrt{w} b(w)$

$$(a(w) + \sqrt{w} b(w)) \left(a'(w) + \frac{1}{2} w^{-\frac{1}{2}} b(w) + \sqrt{w} b'(w) \right) = 1 + a(w) + \sqrt{w} b(w)$$

and compare coefficients of $w^0, w^{\frac{1}{2}}, w^1, w^{\frac{3}{2}}, \dots$

§3. Stationary phase / steepest descent ... method.

Example $I(\lambda) = \int_0^1 \frac{e^{-i\lambda t^2}}{1+t} dt$ as $\lambda \rightarrow \infty$
 (assume $\lambda \in \mathbb{R}_{>0}$).

Issue: $|e^{-i\lambda t^2}| = 1$ for $t \in (0, 1)$.

Consider the function $\phi(t) = it^2$. ($t \in \mathbb{C}$) More traditionally

$$\phi(z) = iz^2 = i(x^2 - y^2 + 2xyi) = -2xy + i(x^2 - y^2)$$

$$|e^{-\lambda\phi(z)}| = e^{-\text{Re}(\lambda\phi(z))} = e^{\lambda(2xy)}$$

→ If we are integrating along a contour $C: z_0 \rightarrow \infty$ along which $2x_0y_0 > 2xy \forall z = x+iy$, then Laplace's method

applies and $\int_C \frac{e^{-i\lambda t^2}}{1+t} dt$ can be computed from the behaviour

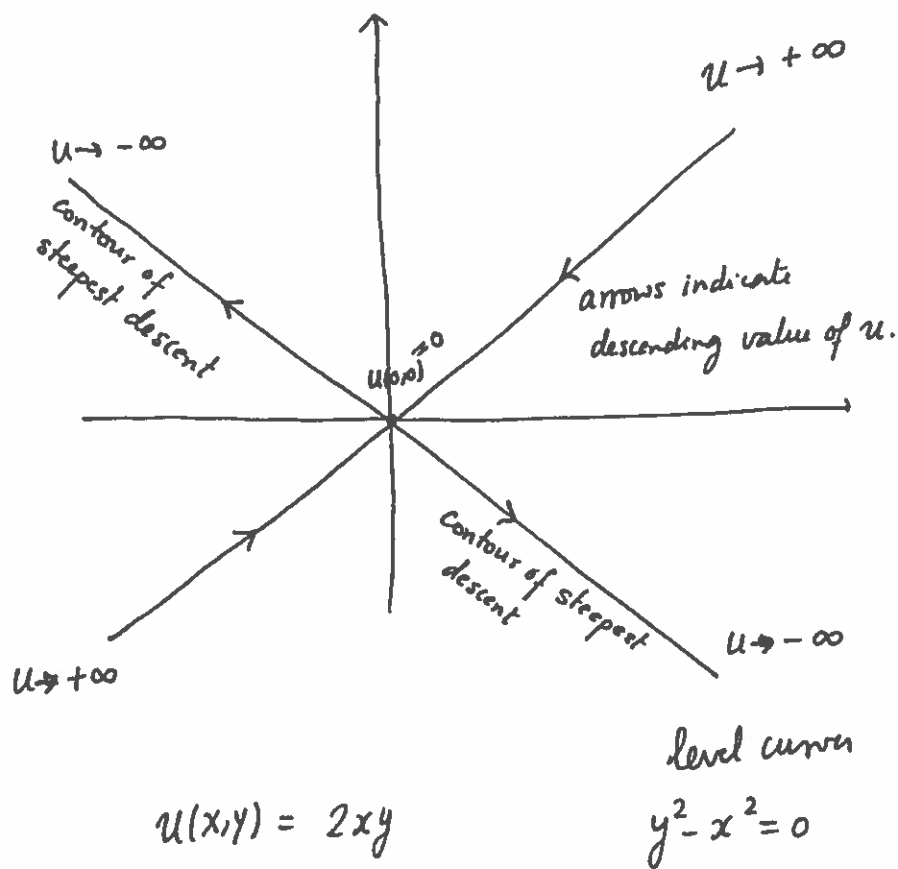
of integrand near z_0 . ["Steepest descent contour"]
 $\psi(z) = -iz^2$

Rewriting $I(\lambda) = \int_0^1 \frac{e^{\lambda(2xy - i(x^2 - y^2))}}{1 + \underbrace{(x+iy)}_z} dz$

For the endpoint 0, $\psi(0) = 0$

Recall: for $\psi(z) = u(x,y) + i v(x,y)$, the curve of steepest ascent/descent for $u(x,y)$ are the level curves of $v(x,y)$.

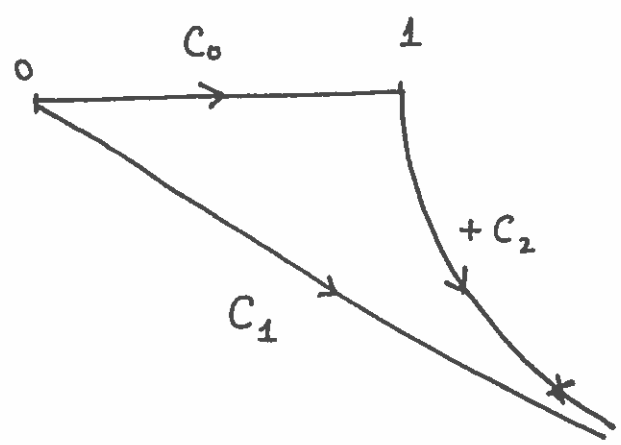
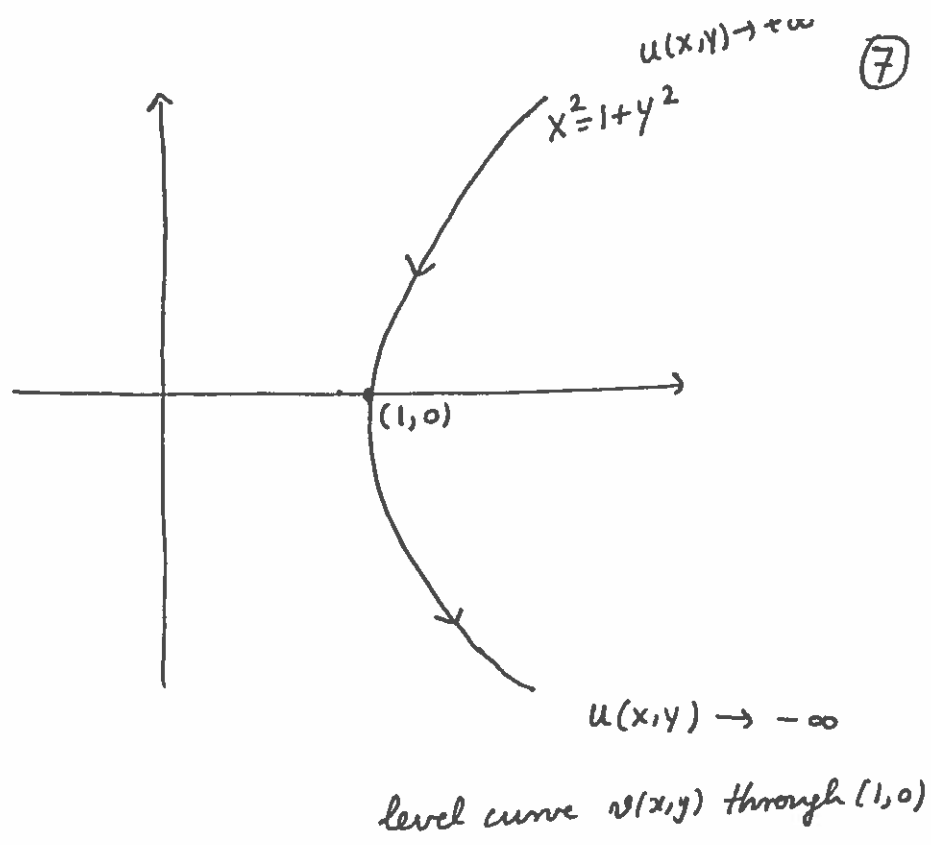
In our case $v(x,y) = y^2 - x^2$. Its level curves through $(0,0)$ are $y = \pm x$



For the endpoint 1, $\psi(1) = -i$

Level curves for $y^2 - x^2 = -1$, $x^2 = 1 + y^2$ are hyperbolas

Now, we argue
 by principle of
 contour deformation
 to replace $[0, 1]$
 by contours of steepest descent.



$$\int_{C_0} = \int_{C_1} - \int_{C_2}$$

Easy exercise - justify this.

Analysis of C_1 = $e^{-i\frac{\pi}{4} \cdot t}$
 $(0 < t < \infty)$

set $z = e^{-i\frac{\pi}{4}} t$, so $z^2 = -it^2$
 $-iz^2 = -t^2$

$$I_1(\lambda) = \int_0^{\infty} \frac{e^{-\lambda t^2}}{1 + e^{-i\frac{\pi}{4}} t} e^{-i\frac{\pi}{4}} dt \quad (\text{back in familiar form})$$

Analysis of C_2

$$I_2(\lambda) = e^{-i\lambda} \int_0^{\infty} \frac{e^{-\lambda p}}{1 + Z(p)} Z'(p) dp \quad \left(\frac{Z'(p)}{1+Z(p)} = \frac{i}{2} (1+ip)^{-1/2} \right)$$

$$(1 + (1+ip)^{1/2})^{-1}$$