

Recall: we were studying the problem of determining $\lambda \rightarrow \infty$ ($\lambda \in \mathbb{R}$)

behaviour of $I(\lambda) = \int_C f(z) e^{\lambda \psi(z)} dz$. Here,

C is a (finite or infinite) contour in $\Omega \subset \mathbb{C}$; $f, \psi: \Omega \rightarrow \mathbb{C}$ are holomorphic.

Note: if $\psi(x+yi) = u(x,y) + i v(x,y)$, then

$|e^{\lambda \psi(z)}| = e^{\lambda \cdot u(x,y)}$. Therefore, if C has an endpoint, say z_0 , so that $u(x_0, y_0) = \text{Max} \{u(x,y) : (x,y) \in C\}$, then

asymptotic expansion of $I(\lambda)$, as $\lambda \rightarrow \infty$, can be obtained via Laplace's method + Watson's lemma.

§1. Definition. Let $C: (a,b) \rightarrow \Omega$ be a path in Ω ; $\psi: \Omega \rightarrow \mathbb{C}$ holomorphic

We say C is a path of ascent (resp. descent)

for $u(x,y) = \text{Re } \psi(x+iy)$, if the composition $(a,b) \xrightarrow{C} \Omega \xrightarrow{u} \mathbb{R}$ is a monotonically increasing (resp. decreasing) function.

C is a path of steepest ascent (resp. descent) if

$\vec{C}'(t) = k \vec{\nabla} u(C(t)) \quad \forall t \in (a,b)$ where C is differentiable,

for some $k > 0$ (resp. $k < 0$).

(i.e. tangent line to C at a point p is in the same (resp. opposite) direction as $\vec{\nabla} u$.)

Prop. (see Lecture 1)

Let $z_0 \in \Omega$ be such that $\psi'(z_0) \neq 0$. Let $\psi(z_0) = A + Bi$.

If C is a path of steepest ascent (or descent) passing through z_0 , then $v(x,y) = B \quad \forall (x,y)$ on C .

Proof. If C is given by $\{(x(t), y(t)) : a \leq t \leq b\}$, then

$$\frac{d}{dt} v(x(t), y(t)) = v_x \cdot x'(t) + v_y \cdot y'(t)$$

$$= k (v_x u_x + v_y u_y) \quad (\langle x'(t), y'(t) \rangle = k \vec{\nabla} u)$$

$$= k (v_x u_x + (-u_x) v_x) \quad (\text{by Cauchy-Riemann equations})$$

$$= 0.$$

□

Example. $\psi(z) = z^2 = x^2 - y^2 + 2xyi$; $\psi'(z) = 2z$.

$$u = x^2 - y^2 ; v = 2xy.$$

If $z_0 \neq 0$, say for instance $z_0 = 1+i$ ($u=0 ; v=2$)

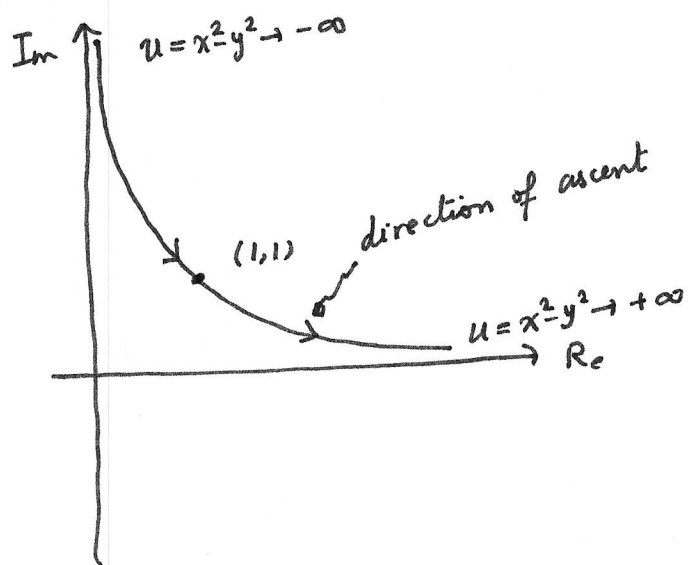
then along $2xy=2$ (or $xy=1$ hyperbola through $(1,1)$)

$x^2 - y^2$ varies fastest from $-\infty$ to ∞

Ex. let $C = \{(\frac{1}{t}, t) : t \geq 1\}$

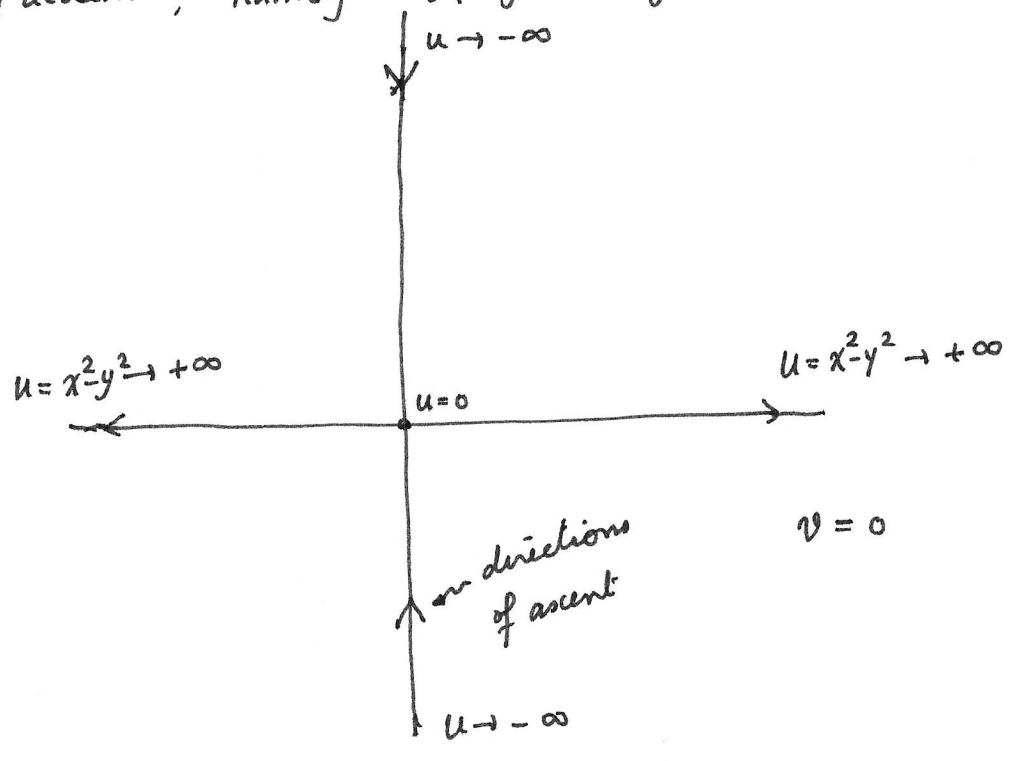
Use Laplace/Watson method to determine the leading behaviour of

$$\int_C \frac{e^{\lambda z^2}}{1+z} dz$$

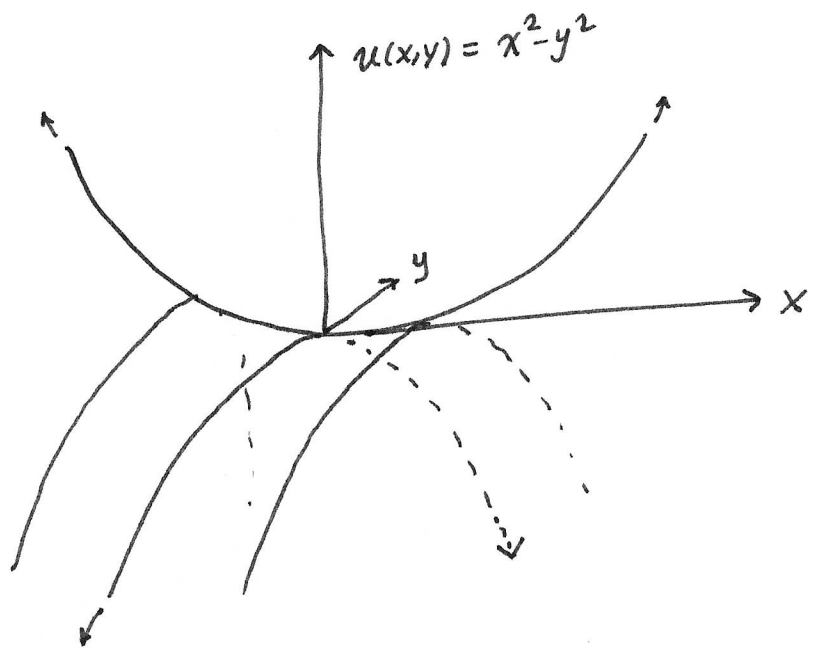


Example contd. ($\psi(z) = z^2$). At $z_0 = 0$, there are 4 directions for steepest ascent/descent, namely $v(x,y) = 2xy = 0$

Thus the graph $\{(x,y, x^2-y^2) : x,y \in \mathbb{R}\}$ has a saddle point at 0



graph of $x^2 - y^2$:



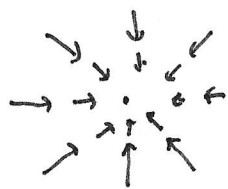
§2. Remarks. - For $\psi(z) = u(x,y) + i v(x,y)$, holomorphic,

Cauchy-Riemann equations imply $\vec{\nabla} u \cdot \vec{\nabla} v = 0$.

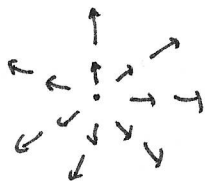
Moreover, $\vec{\nabla} u(x_0, y_0) = 0 \iff \vec{\nabla} v(x_0, y_0) = 0 \iff \psi'(z_0) = 0$.

Defn. z_0 is a critical point of ψ if $\psi'(z_0) = 0$.

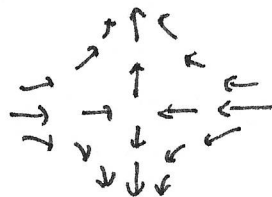
In general - for a \mathbb{R} -valued function $f(x,y)$, $\vec{\nabla} f$ has the following possibilities near a critical point:



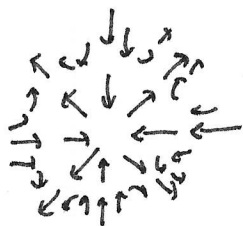
Local Max.



Local Min



Saddle point (order 1)



Saddle point (order 2)

The point of the previous proposition is that for f harmonic, critical points can only be saddle points.

Path of steepest ascent is determined by $x'(t) = u_x(x(t), y(t))$
 $y'(t) = u_y(x(t), y(t))$

Previous proposition solves this system "implicitly" as

$v(x(t), y(t)) = \text{constant}$. (where v is the harmonic conjugate of u)

§3. Method of steepest descent - example (even more basic) ⑤

$$I(\lambda) = \int_0^1 \ln(t) e^{i\lambda t} dt.$$

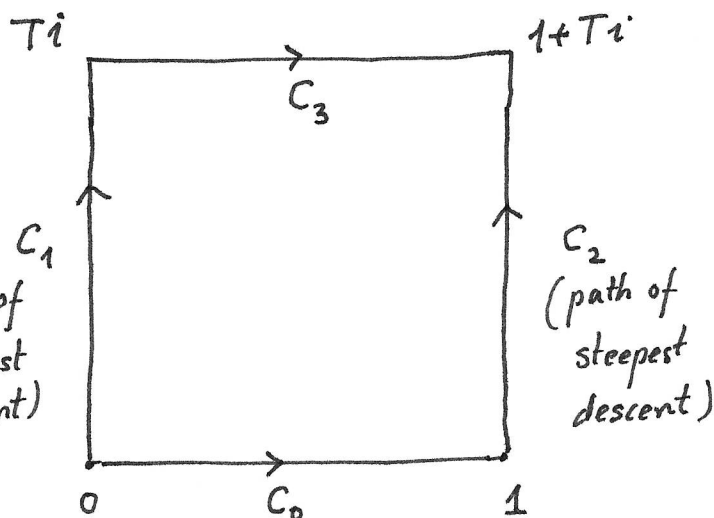
$\psi(z) = iz$, so $u(x,y) = -y$.
 $\psi'(z) = i \neq 0$.
 $v(x,y) = x$

The path of steepest descent of $u(x,y)$ from any (x_0, y_0) is the vertical line (upwards) $x = x_0$.

By Cauchy's theorem,

$$I(\lambda) = \int_{C_1} - \int_{C_2} + \int_{C_3} \log(z) e^{i\lambda z} dz$$

(path of steepest descent)



Claim. $\lim_{T \rightarrow \infty} \int_{Ti}^{1+Ti} \log(z) e^{i\lambda z} dz = 0.$

Pf. $\int_{Ti}^{1+Ti} \log(z) e^{i\lambda z} dz = \int_0^1 \log(u+Ti) e^{i(u+Ti)\lambda} du$
($z = u+Ti$)

$$= e^{-\lambda T} \int_0^1 \log(u+Ti) e^{iu\lambda} du$$

$$\text{Now } \log(u+Ti) = \frac{1}{2} \ln(u^2+T^2) + i \arctan\left(\frac{T}{u}\right) \quad (0 < u < 1)$$

$$\Rightarrow |\log(u+Ti)|^2 < \frac{1}{4} (\ln(1+T^2))^2 + \frac{\pi^2}{4}$$

$$\text{So, } \left| \int_{Ti}^{1+Ti} \log(z) e^{iz\lambda} dz \right| \leq e^{-\lambda T} \left(\frac{1}{4} (\ln(1+T^2))^2 + \frac{\pi^2}{4} \right)^{\frac{1}{2}} \\ \rightarrow 0 \text{ as } T \rightarrow \infty. \quad \square$$

$$\text{So, } I(\lambda) = \lim_{T \rightarrow \infty} \int_0^{Ti} \log(z) e^{iz\lambda} dz - \int_1^{1+Ti} \log(z) e^{iz\lambda} dz$$

Calculation along C_1 ; $z = it$,

$$I_1(\lambda) = \int_0^{\infty} \log(it) e^{-t\lambda} i dt$$

$$= i \int_0^{\infty} \left(\ln(t) + i \frac{\pi}{2} \right) e^{-t\lambda} dt = -\frac{\pi}{2} \frac{1}{\lambda} + i \int_0^{\infty} \ln(t) e^{-t\lambda} dt.$$

$$\int_0^{\infty} \ln(t) e^{-t\lambda} dt = \int_0^{\infty} \ln\left(\frac{u}{\lambda}\right) e^{-\frac{u}{\lambda}} \frac{du}{\lambda} = \frac{1}{\lambda} \left(\int_0^{\infty} \ln(u) e^{-u} du - \ln(\lambda) \int_0^{\infty} e^{-u} du \right) \\ = \frac{-\gamma - \ln(\lambda)}{\lambda}$$

Hence, $I_1(\lambda) = -\frac{\pi}{2\lambda} - i \frac{\gamma + \ln(\lambda)}{\lambda}$ (This is exact!) (7)

Calculation along C_2 . $z = 1 + it$

$$I_2(\lambda) = \int_0^{\infty} \log(1+it) e^{i(1+it)\lambda} i dt$$

$$= i \cdot e^{i\lambda} \int_0^{\infty} \boxed{\log(1+it)} e^{-t\lambda} dt$$

↳ $\sum_{n=1}^{\infty} (-1)^{n-1} \frac{(it)^n}{n}$

By Watson's lemma: $I_2(\lambda) \sim i e^{i\lambda} \left(\sum_{n=1}^{\infty} (-1)^{n-1} i^n \cdot (n-1)! \lambda^{-n-1} \right)$
as $\lambda \rightarrow \infty$.

Hence, $I(\lambda) \sim -\frac{\pi}{2\lambda} - i \frac{\gamma + \ln(\lambda)}{\lambda} - i e^{i\lambda} \left(\sum_{n=1}^{\infty} (-1)^{n-1} i^n (n-1)! \lambda^{-n-1} \right)$

□