

Lecture 21

①

(21.0) Recall : for a smooth connected m -dim'l manifold M , we defined :

- $\Omega^r(M) = \{r\text{-forms on } M\}$ ($0 \leq r \leq m$)
- = smooth sections of $\wedge^r T^*M$

$$(\Omega^0(M) = C^\infty(M))$$

• $\forall \phi \in \Omega^m(M)$, we have $\int_M \phi : C_c(M) \longrightarrow \mathbb{R}$

\mathbb{R} -linear functional \uparrow

continuous $f : M \rightarrow \mathbb{R}$ st. $\text{Supp}(f)$ is compact.

(21.1) de Rham differential (optional)

$$d : \Omega^r(M) \longrightarrow \Omega^{r+1}(M)$$

• Let $M = \bigcup M_\alpha$ be an open covering by coord. open sets

$$M_\alpha \xrightarrow[\substack{\text{homeo.} \\ \sim \\ (x_1^\alpha, \dots, x_m^\alpha)}]{\sim} \text{Cube}_{r_\alpha}(\underline{0}) \subset \mathbb{R}^m$$

$$\omega \in \Omega^r(M) \iff \omega_\alpha = \omega|_{M_\alpha} \quad \forall \alpha \text{ st.}$$

$$\omega_\alpha|_{M_\alpha \cap M_\beta} = \omega_\beta|_{M_\alpha \cap M_\beta}$$

• Locally $\omega_\alpha = \sum_{1 \leq i_1 < \dots < i_r \leq m} f_{\alpha; i_1 \dots i_r} dx_{i_1}^\alpha \wedge \dots \wedge dx_{i_r}^\alpha$

Where for each $\underline{i} = i_1 < \dots < i_r$, $f_{\alpha; \underline{i}}$ is a smooth function on M_α . (2)

• Consistency on overlaps $M_\alpha \cap M_\beta$: recall that the transition (or change of coordinate) functions are assumed to be smooth functions $F_{i; \alpha}^\beta(x_1, \dots, x_m)$ (\mathbb{R} -valued of m real variables)

$$x_i^\beta(q) = F_{i; \alpha}^\beta(x_1^\alpha(q), \dots, x_m^\alpha(q)) \quad \forall q \in M_\alpha \cap M_\beta$$

$$\Rightarrow dx_i^\beta = \sum_{j=1}^m \frac{\partial F_{i; \alpha}^\beta}{\partial x_j^\alpha} \cdot dx_j^\alpha$$

We get (Lecture 20, page 3):

$$dx_{i_1}^\beta \wedge \dots \wedge dx_{i_r}^\beta = \sum_{j_1 < \dots < j_r} \underbrace{\Delta_{\underline{i}, \underline{j}}(J_\alpha^\beta)}_{\uparrow} \cdot dx_{i_1}^\alpha \wedge \dots \wedge dx_{i_r}^\alpha$$

$$\text{set } J_\alpha^\beta = \left[\frac{\partial F_{i; \alpha}^\beta}{\partial x_j^\alpha} \right]_{1 \leq i, j \leq m}$$

$\Delta_{\underline{i}, \underline{j}}(J_\alpha^\beta)$ = determinant of $r \times r$ submatrix of J_α^β
w/ rows $i_1 \dots i_r$ & cols. $j_1 \dots j_r$

So we get (by an easy computation) $\omega_\alpha = \omega_\beta$ on $M_\alpha \cap M_\beta$

iff $\forall (i_1 < \dots < i_r) =: \underline{i}$

$$(*) \quad f_{\alpha; \underline{i}}(x_1^\alpha, \dots, x_m^\alpha) = \sum_{\substack{\underline{l} = \\ (l_1 < \dots < l_r)}} f_{\beta; \underline{l}}(x_{l_1}^\beta, \dots, x_{l_r}^\beta) \cdot \Delta_{\underline{l}, \underline{i}}(J_\alpha^\beta)$$

Thus $\omega \in \Omega^r(M)$ is completely determined by (3)
 $\{f_{\alpha; \underline{i}} \in C^\infty(M_\alpha)\}_{\alpha, i_1 < \dots < i_r}$ satisfying $(*)$, and conversely.

• Define $d\omega \in \Omega^{r+1}(M)$ locally by

$$d\omega|_{M_\alpha} = \sum_{i_1 < \dots < i_{r+1}} \left\{ \sum_{t=1}^{r+1} (-1)^t \frac{\partial f_{\alpha; i_1, \dots, \hat{i}_t, \dots, i_{r+1}}}{\partial x_{i_t}^\alpha} \right\} dx_{i_1}^\alpha \wedge \dots \wedge dx_{i_{r+1}}^\alpha$$

Easy exercise - consistency equation holds $(*)$. Hence we get $d: \Omega^r(M) \rightarrow \Omega^{r+1}(M)$.

Another easy exercise: $d^2 = 0$. de Rham complex is:

$$\Omega^*(M): 0 \rightarrow \Omega^0(M) \xrightarrow{d} \Omega^1(M) \xrightarrow{d} \dots \rightarrow \Omega^m(M) \rightarrow 0$$

de Rham cohomology

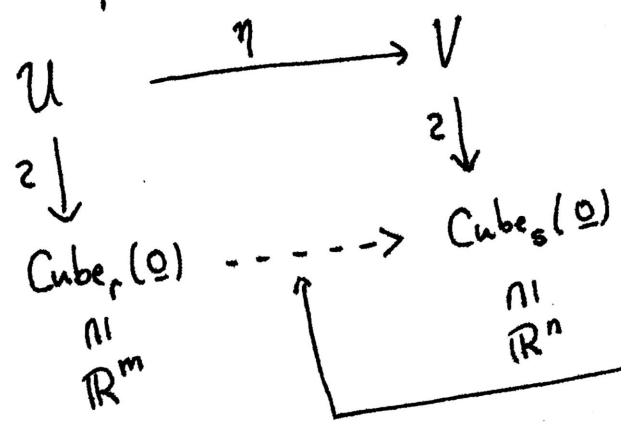
$$H_{dR}^i(M) := \frac{\text{Ker}(\Omega^i(M) \rightarrow \Omega^{i+1}(M))}{\text{Im}(\Omega^{i-1}(M) \rightarrow \Omega^i(M))} \quad (0 \leq i \leq m)$$

• Compute H^0 and H^1 of S^1 as an ~~is~~ instructive exercise.

(21.2) Pull back of differential forms.

(4)

Let $\eta: M \rightarrow N$ be a smooth map. Recall that η being smooth means $\forall p \in M$, we can find a coordinate neighbourhood of $q = \eta(p) \in N$ and $p \in M$, say $U \subset M$, $V \subset N$ s.t. $\eta^{-1}(V) = U$ and



is smooth fn.

Thus we can cover $M = \cup M_\alpha$ and $N = \cup N_\alpha$ s.t. $\eta(M_\alpha) \subset N_\alpha \subset N$. This allows us to define the pull-back of $\omega \in \Omega^r(N)$, denoted by $\eta^* \omega \in \Omega^r(M)$ locally via:

$$\omega|_{N_\alpha} = \sum_{1 \leq i_1 < \dots < i_r \leq n} f_{\alpha; i_1 \dots i_r} dx_{i_1}^\alpha \wedge \dots \wedge dx_{i_r}^\alpha$$

$$\text{then } \eta^* \omega|_{M_\alpha} := \sum_{\underline{i}} (f_{\alpha; \underline{i}} \circ \eta) d(x_{i_1}^\alpha \circ \eta) \wedge \dots \wedge d(x_{i_r}^\alpha \circ \eta)$$

$$\text{where } M_\alpha \xrightarrow{\eta} N_\alpha \xrightarrow{(x_1^\alpha, \dots, x_n^\alpha)} \text{Cube}_{S_\alpha}(\mathbf{0}) \subset \mathbb{R}^n$$

One need to check (easy!) that $\{\eta^* \omega|_{M_\alpha}\}$ satisfy (*) of page 2 above.

(21.3) Back to Lie groups. Let G be a connected Lie group of dimension n . An r -form $\omega \in \Omega^r(G)$ is said to be left-invariant if $\forall \sigma \in G$ $l_\sigma^* \omega = \omega$. Here $l_\sigma : G \rightarrow G$ is left mult. $x \mapsto \sigma x$ (5)

(Prop (14) Lecture 1 page 5) $\mathfrak{g} = T_e G =$ left-invariant vector fields on G

Analogously, $\mathfrak{g}^* =$ left-invariant 1-forms on G . If we pick a basis x_1, \dots, x_n of \mathfrak{g} , we obtain dx_1, \dots, dx_n (dual basis of \mathfrak{g}^*) 1-forms on G .

$\wedge^n \mathfrak{g}^*$ is 1-dim'l vector space (over \mathbb{R}). Thus up to a scalar $\phi = dx_1 \wedge \dots \wedge dx_n$ is the only left-invariant n -form on G .

Note, if $\phi(\sigma) = 0$ for some $\sigma \in G$, then $\phi(\tau) = 0 \forall \tau \in G$ by invariance. Since we know $\phi \neq 0$ element of $\wedge^n \mathfrak{g}^*$, we get that ϕ is nowhere vanishing, hence G is orientable.

From now on, we fix ϕ and assume G is oriented via ϕ (see Lecture 20 page 4).

(21.4) Modularity. For $\sigma_0 \in G$, consider $\text{Conj}(\sigma_0) : G \rightarrow G$
 $x \mapsto \sigma_0 x \sigma_0^{-1}$

Then $\text{Conj}(\sigma_0)^* \phi$ is another left-invariant n -form on G . \textcircled{c}

$$\Rightarrow \exists c(\sigma_0) \in \mathbb{R}_{\neq 0} \text{ s.t. } \text{Conj}(\sigma_0)^* \phi = c(\sigma_0) \cdot \phi.$$

In terms of \int_G defined using ϕ , we get the following

identities

$$(i) \int_G (f \circ l_\sigma) \phi = \int_G f \cdot \phi \quad \forall \sigma \in G$$

Let $r_{\sigma_0} : G \rightarrow G$ be the right multiplication. Then

$$\text{Conj}_{\sigma_0}^{-1} = l_{\sigma_0}^{-1} \circ r_{\sigma_0} \text{ by definition. And we get}$$

$$\begin{aligned} (ii) \int_G (f \circ r_{\sigma_0}) \phi &= \int_G (f \circ l_{\sigma_0} \circ \text{Conj}(\sigma_0)^{-1}) \phi \\ &= \int_G (f \circ l_{\sigma_0}) (\text{Conj}(\sigma_0)^* \phi) = c(\sigma_0) \int_G (f \circ l_{\sigma_0}) \phi \\ &= c(\sigma_0) \int_G f \cdot \phi. \end{aligned}$$

(21.5) Let G be a compact (connected) Lie group. We normalize (7)

ϕ so that $\int_G 1 \cdot \phi = 1$ constant function 1.

From (ic) of page 6, we get $c(\sigma_0) = 1 \quad \forall \sigma_0 \in G$.

We fix this n -form (both left & right invariant) and will use the familiar notation $\int_G f(x) dx$ for $\int f \cdot \phi$.

Note: Lie groups for which $c(\sigma_0) = 1 \quad \forall \sigma_0$ are called unimodular. All compact Lie groups are unimodular.

Similar to the computation (ii) of page 6, gives $\int_G f(\bar{x}') dx = \int_G f(x) dx$

(21.6) Consequences of $\int_G f(x) dx$:

(a) Let V be a finite-dimensional repr. of G , over \mathbb{C} .

That is, V is a f.d. vector space over \mathbb{C} and we have

a continuous map $G \longrightarrow GL(V)$.

(group hom.)

Then V admits a G -invariant Hermitian form ⑧

Pf. Let $\langle \cdot, \cdot \rangle : V \times V \rightarrow \mathbb{C}$ be any Hermitian form.

Define $H(v, w) = \int_G \langle g.v, g.w \rangle dg$. Then $H(\cdot, \cdot)$ is

the required G -invariant Hermitian form. □

(b) If $G \curvearrowright V$ (f.d. repr. / \mathbb{C}) and $V' \subset V$ is a subrepr.

Then $\exists V'' \subset V$ a subrepr. s.t. $V = V' \oplus V''$.

(take $V'' = (V')^\perp$ w.r.t. a G -invariant Hermitian form).

(c) Schur's Lemma. V_1, V_2 f.d. irreducible reprs. of G . Then

$$\text{Hom}_G(V_1, V_2) = \begin{cases} 0 & \text{if } V_1 \not\cong V_2 \\ \mathbb{C} & \text{if } V_1 \cong V_2 \end{cases}$$

Pf. If $\xi : V_1 \rightarrow V_2$ is a hom. of G -reprs. then

$\text{Ker}(\xi) \subset V_1$ is a subrepr. $\Rightarrow \begin{cases} \text{Ker}(\xi) = V_1 \Rightarrow \xi = 0 \\ \text{Ker}(\xi) = 0 \Rightarrow V_1 \subset V_2 \text{ is a} \\ \text{subrepr. By irr. of } V_2 \end{cases}$

$\xi : V_1 \rightarrow V_2$ is iso.

In the latter case, let $\lambda \in \mathbb{C}$ be an

eigenvalue of ξ . Then $\text{Ker}(\xi - \text{Id} \cdot \lambda) \subset V_1$ is a ~~non-zero~~ non-zero

subrepr. $\Rightarrow \text{Ker}(\xi - \lambda \cdot \text{Id}) = V \Rightarrow \xi = \lambda \cdot \text{Id}$. □ (9)

(d) Schur's orthogonality relations. Let V_1, V_2 be two irr. f.d. reprs. of G , with G -invariant Hermitian forms $H_1(\cdot, \cdot)$ and $H_2(\cdot, \cdot)$ resp.

(i) $V_1 \not\cong V_2$. For every $u_1, v_1 \in V_1$ and $u_2, v_2 \in V_2$ we have

$$\int_G H_1(gu_1, v_1) \overline{H_2(gu_2, v_2)} dg = 0$$

(ii) If $V_1 = V_2$, $\int_G (gu_1, v_1) \overline{(gu_2, v_2)} dg = \frac{(u_1, u_2) \overline{(v_1, v_2)}}{\dim V}$

pp for any \mathbb{C} -linear map $\xi: V_2 \rightarrow V_1$, let $L: V_2 \rightarrow V_1$ be defined by $\forall v_2 \in V_2$ $L(v_2) = \int_G g \cdot (\xi(\bar{g}^{-1} \cdot v_2)) dg \in V_1$. Then

$L: V_2 \rightarrow V_1$ is a hom. of reprs. of G . Thus $L \equiv 0$ if $V_1 \not\cong V_2$

in which case: $0 = H_1(Lv_2, v_1) = \int_G (g \xi \bar{g}^{-1} \cdot v_2, v_1) dg$

(take $\xi(v_2') = (v_2', u_2) \cdot u_1$)

$$= \int_G (g (\bar{g}^{-1} v_2, u_2) u_1, v_1) dg$$

$$= \int_G (gu_1, v_1) \overline{(gu_2, v_2)} dg$$

Now, if $V_1 = V_2$, we get $L = \lambda \cdot \text{Id}$ where $\lambda \cdot \dim(V)$ (10)

$= \text{Tr}(L) = \text{Tr}(\xi)$. Again take $\xi(w) = (w, u_2) \cdot u_1$ to get

$$\left(\frac{\text{Tr}(\xi)}{\dim V} \right)_{(V_1, V_2)} = \frac{(u_1, u_2)}{\dim V} \overline{(V_1, V_2)}$$
$$\downarrow = \lambda \overline{(V_1, V_2)} = (L v_2, v_1) = \int_G (g \xi \bar{g}^{-1} \cdot v_2, v_1) dg$$

$$= \int_G (\bar{g}^{-1} v_2, u_2) (g u_1, v_1) dg$$

$$= \int_G (g u_1, v_1) \overline{(g u_2, v_2)} dg \quad \text{as required} \quad \square$$