

## Lecture 4.

(1)

Recall: Laplace transform along direction  $\theta$

$$(\mathcal{L}_\theta \phi)(z) = \int_{l_\theta} \phi(t) e^{-zt} dt$$

$$l_\theta = \mathbb{R}_{>0} e^{i\theta}$$

Assuming  $\phi: \mathbb{C} \dashrightarrow \mathbb{C}$  is meromorphic with discrete set of poles  $A \subset \mathbb{C}$  ( $0 \notin A$ ), and sub-exponential growth at  $\infty$  (so we can take its Laplace transform)

$(\mathcal{L}_\theta \phi)(z)$  makes sense for every  $\theta$  s.t.  $\mathbb{R}_{>0} e^{i\theta} \cap A = \emptyset$ .

§1. Some definitions .. We say  $l_\theta$  is a Stokes ray if  $A \cap l_\theta \neq \emptyset$ .  
(for  $\phi$ )

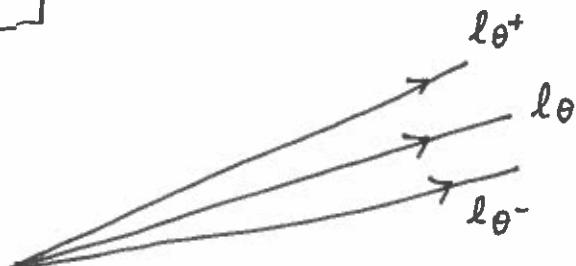
Stokes' automorphism  $S_\theta$  is defined by the following formula

$$\boxed{\mathcal{L}_{\theta^+} \phi = \mathcal{L}_{\theta^-} (S_\theta \phi)}$$

where  $l_\theta^\pm$  are slightly above/below

$l_\theta$  s.t.  $A \cap l_\theta^\pm = \emptyset$ .

(see picture)



(2)

Full discontinuity across  $\ell_\theta$  is defined by

$$S_\theta = \text{Id} - \text{Disc}_\theta, \text{ so we have}$$

$$\mathcal{L}_{\theta^+} = \mathcal{L}_{\theta^-} \circ (\text{Id} - \text{Disc}_\theta) \text{ and hence } \mathcal{L}_{\theta^+} - \mathcal{L}_{\theta^-} = -\mathcal{L}_{\theta^-} \circ \text{Disc}_\theta.$$

The logarithm of Stokes' automorphism defines "alien derivatives"

(Écalle - derivation étranger)

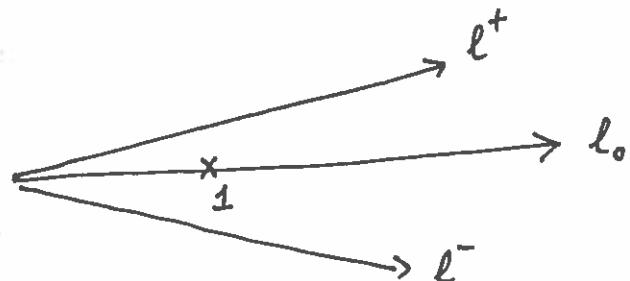
$$\boxed{S_\theta = \exp \left( \sum_{\omega \in A_\theta = A \cap \ell_\theta} e^{-\omega z} \Delta_\omega \right)} - (*) \begin{pmatrix} \text{see §3 below} \\ \text{for some} \\ \text{remarks} \end{pmatrix}$$

§2. Example 1.  $\phi(t) = \frac{1}{1-t}$  ( $=$  Borel transform of  $f(z) = \sum_{n=0}^{\infty} n! z^{-n-1}$ )

$$S_\theta \phi = \phi \text{ for } \theta \neq 0 \bmod 2\pi$$

$$(\mathcal{L}_+ \phi - \mathcal{L}_- \phi)(z)$$

$$= 2\pi i e^{-z}$$



Note - both  $(\mathcal{L}_\pm \phi)(z)$  solve

$$F'(z) + F(z) = \bar{z}^1$$

and  $e^{-z}$  is a soln. of the associated hgs. equation  $F' + F = 0$ .

(3)

Formula (\*) , in this case, can be written as

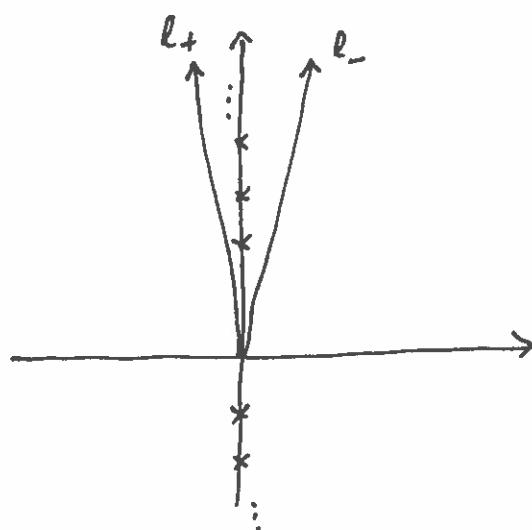
$$\mathcal{L}_{\theta^+} \phi = \mathcal{L}_{\theta^-} \phi + \sum_{k=1}^{\infty} e^{-kz} \mathcal{L}_{\theta^-} (\Delta_1^k \cdot \phi)$$

Comparing with  $(\mathcal{L}_{\theta^+} - \mathcal{L}_{\theta^-})(\phi) = 2\pi i e^{-z}$ , we get

$$\Delta_1 \left( \frac{1}{1-t} \right) = 2\pi i \quad \Delta_1^n \left( \frac{1}{1-t} \right) = 0 \quad \forall n \geq 2.$$

Example 2.  $\phi(t) = \frac{t}{1-e^{-t}}$

$$A = 2\pi i \mathbb{Z}_{\neq 0}$$



$$(\mathcal{L}_+ \phi)(z) - (\mathcal{L}_- \phi)z$$

$$= -(2\pi i)^2 \sum_{l=1}^{\infty} l \cdot e^{-2\pi i l z}$$

$$= 4\pi^2 \sum_{l=1}^{\infty} l \cdot e^{-2\pi i l z}$$

$$\text{Check. } - \quad \Delta_{2\pi i l} \cdot \phi = 4\pi^2 l \quad \forall l \geq 1.$$

§3. Some remarks on the definition and generalizations.

(i) The definition given above is a bit imprecise. - since

terms of the form  $e^{-\omega z}$  are asymptotically 0 and as such

(4)

cannot be obtained from asymptotic power series analysis alone. It is best to view  $\bar{e}^{-\omega z}$ -type terms symbolically thus defining  $\Delta_\omega \phi$  as a function whose Laplace transform gives the coefficient of  $\bar{e}^{-\omega z}$  in  $\log(S_\theta)$  - see (\*) on page 2 above.

(ii) Even with these caveats,  $(\mathcal{L}_0 \phi)(z) = c \in \mathbb{C}$  does not hold

for any  $\phi$ . Recall - originally -  $\bar{z}' \mathbb{C}[[\bar{z}']] \xrightarrow{\mathcal{B}} \mathbb{C}[[t]]$

$$\begin{array}{ccc} & \uparrow & \\ & \mathcal{B} & \\ & \downarrow & \\ z' & & \end{array}$$

no constant term.

The convention employed in the literature is to introduce a

symbol  $\delta$  s.t.  $\mathcal{B}(c) = c\delta$  ( $\mathcal{L}\delta = 1$ )

( $\delta$  = unit for convolution product - to be discussed later)

$$\mathbb{C} + \bar{z}' \mathbb{C}[[\bar{z}']] \xrightarrow{\sim} \mathbb{C} \cdot \delta + \mathbb{C}[[t]]$$

Hence, the precise answer from Examples 1, 2 should be

$$\Delta_1 \left( \frac{1}{1-t} \right) = 2\pi i \delta$$

$$\Delta_{2\pi i l} \left( \frac{t}{1-\bar{e}^{-t}} \right) = 4\pi^2 l \delta \quad (l \geq 1)$$

§4. Allowing for log-type singularities in the Borel plane. - review of gamma function.-

$$\int_0^\infty t^{x-1} e^{-t} dt = \Gamma(x) \quad \text{for } \operatorname{Re}(x) > 0. \quad (**)$$

is the original Euler's definition of the gamma function.

$$\text{Here, } t^{x-1} = e^{(x-1)\ln(t)} \quad (t \in \mathbb{R}_{>0}).$$

Exercise. - Show that

$$\int_0^\infty t^a e^{-tz} dt = \Gamma(a+1) z^{-a-1} \quad \begin{array}{l} \operatorname{Re}(z) > 0 \\ \operatorname{Re}(a) > -1. \end{array}$$

Justify differentiation <sup>(w.r.t. a)</sup> under the integral sign to get

$$\int_0^\infty \ln(t) t^a e^{-tz} dt = \Gamma'(a+1) z^{-a-1} + \frac{1}{(a+1)z^2} \Gamma(a+1) z^{-a-1} - \log(z) \Gamma(a+1) z^{-a-1}$$

Set  $a=0$  to get

$$\int_0^\infty \ln(t) e^{-tz} dt = (\Gamma'(1) - \log(z)) z^{-a-1}.$$

$\Gamma(1)=1$  and  $\Gamma'(1)=-\gamma$ .

(6)

Thus Laplace transform is defined for more general class of functions (e.g.  $t^a = e^{a \log(t)}$ : "multi-valued")

The previous exercise allows us to generalize Watson's lemma (see Lecture 3, § 2).

$$\left[ \begin{array}{l} \text{If } \phi(t) \sim \sum_{n=1}^{\infty} c_n t^{a_n} \quad -1 < a_1 < a_2 < \dots \\ \qquad \qquad \qquad \text{as } t \rightarrow 0^+ \qquad \qquad \text{real numbers} \\ \text{then} \quad \int_0^{\infty} \phi(t) e^{-zt} dt \sim \sum_{n=0}^{\infty} c_n \Gamma(a_{n+1}) z^{-a_{n+1}} \quad \text{as} \\ \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \text{Re}(z) \rightarrow \infty \end{array} \right]$$