

## Lecture 25

(1)

### Linear integral equations - Fredholm theory

§1. Integral equations were first introduced by Laplace in 1782 who considered the following integral transformations:

$$f(x) = \int e^{xt} \phi(t) dt \quad g(x) = \int t^{x-1} \phi(t) dt$$

Next significant appearance of integral equations was in the works of Fourier (1802)

$$f(x) = \int_{-\infty}^{\infty} \cos(xt) \phi(t) dt \quad \leadsto \quad \phi(x) = \frac{2}{\pi} \int_0^{\infty} \cos(ux) f(u) du$$

and Abel (1823)

$$f(x) = \int_a^x \frac{u(t)}{(x-t)^\mu} dt \quad 0 < \mu < 1, \quad a \leq x \leq b$$

$$\leadsto \quad u(z) = \frac{\sin \mu \pi}{\pi} \frac{d}{dz} \int_a^z \frac{f(x)}{(z-x)^{1-\mu}} dx.$$

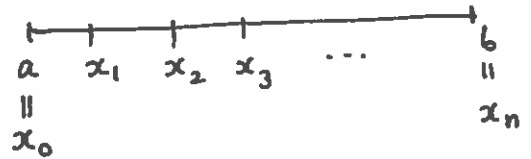
A major step towards a systematic study of such equations was taken by Fredholm in 1903. Fredholm's work has been vastly generalized to "spectral theory of compact operators" - by results and theories of Hilbert spaces.

§2. Fredholm equations. - Given a continuous function  $f(x)$  on  $a \leq x \leq b$  and  $K(x,y)$ ,  $a \leq x,y \leq b$ , solve for  $\phi(x)$  s.t.

$$\phi(x) - \lambda \int_a^b K(x,y) \phi(y) dy = f(x) \quad - (*)$$

Fredholm's derivation of  $\phi(x)$ . -

Consider the subdivision of  $[a,b]$  into  $n$  equal subintervals



$$x_0 = a, \quad x_{p+1} - x_p = \delta \left( = \frac{b-a}{n} \right)$$

$$x_n = b.$$

Writing the partial Riemann sums, and taking (\*) for  $x = x_p$  gives a linear system of equations:

$$\phi(x_p) - \lambda \delta \sum_{q=1}^n K(x_p, x_q) \phi(x_q) = f(x_p) \quad 1 \leq p \leq n$$

By linear algebra, this system is solved by

$$\phi(x_\mu) = \frac{1}{D_n(\lambda)} \cdot \sum_{p=1}^n D_n(x_\mu, x_p) f(x_p), \quad \text{where}$$

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$$D_n(\lambda) = \det \left( \text{Id} - \lambda \delta \left( K(x_p, x_q) \right)_{1 \leq p, q \leq n} \right)$$

$$= 1 - \lambda \delta \sum_{p=1}^n K(x_p, x_p) + \frac{\lambda^2 \delta^2}{2!} \sum_{p, q=1}^n \begin{vmatrix} K(x_p, x_p) & K(x_p, x_q) \\ K(x_q, x_p) & K(x_q, x_q) \end{vmatrix} \\ + \dots + (-1)^n \frac{\lambda^n \delta^n}{n!} \det \left( K(x_p, x_q) \right)_{1 \leq p, q \leq n}.$$

and  $D_n(x_\mu, x_\nu)$  is the cofactor in  $D_n(\lambda)$  of the term involving  $K(x_\nu, x_\mu)$ . In the limit,  $n \rightarrow \infty$  (hence  $\delta \rightarrow 0$ ) we find

$$D(\lambda) = 1 - \lambda \int_a^b K(\xi_1, \xi_1) d\xi_1 + \frac{\lambda^2}{2!} \int_a^b \int_a^b \begin{vmatrix} K(\xi_1, \xi_1) & K(\xi_1, \xi_2) \\ K(\xi_2, \xi_1) & K(\xi_2, \xi_2) \end{vmatrix} d\xi_1 d\xi_2 \\ - \dots$$

and

$$D(x_\mu, x_\nu; \lambda) = \lambda K(x_\mu, x_\nu) - \lambda^2 \int_a^b \begin{vmatrix} K(x_\mu, x_\nu) & K(x_\mu, \xi_1) \\ K(\xi_1, x_\nu) & K(\xi_1, \xi_1) \end{vmatrix} d\xi_1 \\ + \frac{\lambda^3}{3!} \int_a^b \int_a^b \begin{vmatrix} K(x_\mu, x_\nu) & K(x_\mu, \xi_1) & K(x_\mu, \xi_2) \\ K(\xi_1, x_\nu) & K(\xi_1, \xi_1) & K(\xi_1, \xi_2) \\ K(\xi_2, x_\nu) & K(\xi_2, \xi_1) & K(\xi_2, \xi_2) \end{vmatrix} d\xi_1 d\xi_2 \dots$$

Hence our solution should be of the form

$$\phi(x) = f(x) + \frac{1}{D(\lambda)} \int_a^b D(x, \xi; \lambda) f(\xi) d\xi$$

§3. Given  $K(x,y)$  continuous on  $[a,b] \times [a,b]$ , consider

$$D(\lambda) := 1 + \sum_{n=1}^{\infty} a_n \frac{\lambda^n}{n!} \quad ; \quad a_n = (-1)^n \int_{[a,b]^n} \left| \left( K(\xi_i, \xi_j) \right)_{1 \leq i, j \leq n} \right| d\xi_1 \dots d\xi_n$$

$$D(x,y; \lambda) = \sum_{n=0}^{\infty} \frac{v_n(x,y)}{n!} \lambda^{n+1} \quad ; \quad v_n(x,y) = (-1)^n \int_{[a,b]^n} K(\xi_i, \xi_j)_{1 \leq i, j \leq n}$$

where

$$v_n(x,y) = (-1)^n \underbrace{\int_a^b \dots \int_a^b}_{n\text{-fold}} \left| \begin{array}{c} K(x,y) \quad K(x,\xi_1) \dots K(x,\xi_n) \\ K(\xi_1,y) \quad \boxed{K(\xi_i, \xi_j)} \\ \vdots \\ K(\xi_n,y) \end{array} \right| d\xi_1 \dots d\xi_n$$

Hadamard's lemma [If  $|a_{ij}| < M \quad \forall i,j \in \{1, \dots, n\}$ , then  $|\det(a_{ij})| \leq n^{n/2} \cdot M^n$ .

(row sums are bounded as  $\sum_{j=1}^n \frac{a_{ij}^2}{nM^2} \leq 1$ , so  $|\det(\frac{a_{ij}}{\sqrt{n} \cdot M})| \leq 1$ .)  
↑  
(Hadamard's lemma)]

easily implies that  $D(\lambda)$  and  $D(x,y; \lambda)$  are entire functions of  $\lambda$ .

Theorem. - 
$$D(x,y; \lambda) = \lambda D(\lambda) K(x,y) + \lambda \int_a^b D(x, \xi; \lambda) K(\xi, y) d\xi$$

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Proof.- Recall  $D(x, y; \lambda) = \sum_{n=0}^{\infty} v_n(x, y) \frac{\lambda^{n+1}}{n!}$ ,

where  $v_n(x, y) = (-1)^n \int_a^b \dots \int_a^b \begin{vmatrix} K(x, y) & K(x, \xi_1) & \dots & K(x, \xi_n) \\ K(\xi_1, y) & \boxed{K(\xi_i, \xi_j)} & & \\ \vdots & & & \\ K(\xi_n, y) & & & \end{vmatrix} d\xi_1 \dots d\xi_n$

$\underbrace{\int_a^b \dots \int_a^b}_{n\text{-fold}}$

Expanding the determinant along the first column, we get-

Integrand of  $v_n(x, y) = K(x, y) \cdot \left| (K(\xi_i, \xi_j))_{1 \leq i, j \leq n} \right|$   
 $+ \sum_{j=1}^n (-1)^j K(\xi_j, y) \left| \begin{matrix} K(x, \xi_1) & \dots & K(x, \xi_n) \\ \boxed{K(\xi_a, \xi_b)} \\ \vdots \\ K(\xi_n, \xi_j) \end{matrix} \right|$

↑  
renaming variables - we get  
the same answer upon  $\int_{[a, b]^n} (-) d\xi$ .

$\Rightarrow D(x, y; \lambda) = \lambda D(\lambda) K(x, y)$

$+ \sum_{n=1}^{\infty} (-1)^n \frac{\lambda^{n+1}}{n!} \cdot (-1) \cdot n \cdot \underbrace{\int_a^b \dots \int_a^b}_{n\text{ fold}} \begin{vmatrix} K(x, \xi_0) & \dots & K(x, \xi_n) \\ K(\xi_1, \xi) & \dots & K(\xi_n, \xi) \\ \vdots & & \vdots \\ K(\xi_{n-1}, \xi) & \dots & \end{vmatrix} \cdot K(\xi, y) d\xi d\xi_1 \dots d\xi_{n-1}$

$= \lambda D(\lambda) K(x, y) + \lambda \int_a^b D(x, \xi; \lambda) K(\xi, y) d\xi$  as claimed  $\square$

§4. Using the calculation given in the previous section,

the integral eq<sup>n</sup>  $\phi(x) = f(x) + \lambda \int_a^b K(x,y)\phi(y) dy$

gives  ~~$\int_a^b f(x)\mathcal{D}(x)$~~  (multiply on the left by  $\mathcal{D}(x, \xi; \lambda)$  and integrate w.r.t.  $\xi$ )

$$\int_a^b \mathcal{D}(x, \xi; \lambda) f(\xi) d\xi = \int_a^b \mathcal{D}(x, \xi; \lambda) \phi(\xi) d\xi$$

$$- \lambda \int_a^b \int_a^b \mathcal{D}(x, \xi; \lambda) K(\xi, y) \phi(y) dy d\xi$$

Substitute  $\lambda \int_a^b \mathcal{D}(x, \xi) K(\xi, y) d\xi = \mathcal{D}(x, y) - \lambda \mathcal{D}(\lambda) K(x, y)$

to get  $\int_a^b \mathcal{D}(x, \xi) f(\xi) d\xi = \lambda \mathcal{D}(\lambda) \int_a^b K(x, y) \phi(y) dy$   
 $= \mathcal{D}(\lambda) (\phi(x) - f(x))$

So, if  $\mathcal{D}(\lambda) \neq 0$ , Fredholm equation has a unique solution

given by  $\phi(x) = f(x) + \frac{1}{\mathcal{D}(\lambda)} \int_a^b \mathcal{D}(x, \xi) f(\xi) d\xi$ .  $\square$

§5. Some Examples.-

(1) Let  $K(x,y) = xy$  ;  $[a,b] = [0,1]$  .  $f(x) = x$ .

$$D(\lambda) = 1 - \frac{1}{3} \lambda \quad \left( \text{recall: } D(\lambda) = 1 + \sum (-1)^n \frac{a_n \lambda^n}{n!} \right)$$

$$a_n = \int_{[a,b]^n} \det (K(\xi_i, \xi_j)) d\xi$$

eg.  $n=2$

$$\begin{vmatrix} \xi_1^2 & \xi_1 \xi_2 & \xi_1 \xi_3 \\ \xi_2 \xi_1 & \xi_2^2 & \xi_2 \xi_3 \\ \xi_3 \xi_1 & \xi_3 \xi_2 & \xi_3^2 \end{vmatrix} = 0 \quad \text{since all rows are multiple of } (\xi_1 \ \xi_2 \ \xi_3) .$$

$$D(x,y; \lambda) = \lambda xy \quad \text{and} \quad \phi(x) = \frac{3x}{3-\lambda} .$$

$$(2) \quad K(x,y) = xy + y^2 \rightsquigarrow D(\lambda) = 1 - \frac{2}{3} \lambda - \frac{1}{72} \lambda^2$$

$$D(x,y; \lambda) = \lambda(xy + y^2) + \lambda^2 \left( \frac{1}{2} xy^2 - \frac{1}{3} xy - \frac{1}{3} y^2 + \frac{1}{4} y \right)$$

(3) Exercise - from the proof of Thm §3 above, show that:

$$\int_a^b D(\xi, \xi; \lambda) d\xi = -\lambda \frac{d(D(\lambda))}{d\lambda}$$

(4) If  $K(x,y) = f_1(x) f_2(y)$  and  $A = \int_a^b f_1(x) f_2(x) dx$  , then

$$D(\lambda) = 1 - A \lambda \quad \text{and} \quad D(x,y; \lambda) = \lambda f_1(x) f_2(y) .$$

and soln. to Fredholm eq<sup>n</sup>  $\phi(x) - \lambda \int_a^b K(x,y) \phi(y) dy = f(x)$  (8)

is given by  $f(x) + \frac{\lambda f_1(x)}{1 - A\lambda} \int_a^b f(\xi) f_2(\xi) d\xi$ .

Ex. - More generally if  $K(x,y) = \sum_{\ell=1}^m f_{\ell}(x) g_{\ell}(y)$ , then

$D(\lambda)$  and  $D(x,y;\lambda)$  are polynomials in  $\lambda$ .

§6. If  $D(\lambda)$  has a zero at  $\lambda = \lambda_0$  to order  $m$ , i.e.,  
 $D(\lambda) = c_m (\lambda - \lambda_0)^m + \dots$  ( $c_m \neq 0$ ) near  $\lambda = \lambda_0$ ,

and  $D(x,y;\lambda) = g_{\ell}(x,y) (\lambda - \lambda_0)^{\ell} + \dots$ ,

then, by  $\int_a^b D(\xi,\xi;\lambda) d\xi = -\lambda D'(\lambda)$  we see that  
 $m-1 \geq \ell$ ,

Hence  $D(x,y;\lambda) = \lambda D(\lambda) K(x,y) + \lambda \int_a^b K(x,\xi) D(\xi,y) d\xi$

$\Rightarrow g_{\ell}(x,y) = \lambda_0 \int_a^b K(x,\xi) g_{\ell}(\xi,y) d\xi$ .

So, for each  $y$ ,  $g_{\ell}(x,y)$  as a fn. of  $x$ , is a soln. to the

hgs integral eq<sup>n</sup>

$$\psi(x) = \lambda_0 \int_a^b K(x,\xi) \psi(\xi) d\xi$$



§7. Summary of results. -

(1) In case of symmetric kernel,  $K(x,y) = K(y,x)$ , a theorem of Schmidt states that  $D(\lambda) = 0$  has at least one root  $\lambda = \lambda_0$ .

(2) At a root  $\lambda_0$  of  $D(\lambda) = 0$ , if we have  $n$  independent solutions to  $\psi(x) = \lambda_0 \int_a^b K(x,y)\psi(y) dy$ , then

$$n \leq \lambda_0^2 \int_a^b \int_a^b (K(x,y))^2 dx dy . \text{ Hence, the dimension}$$

of the space of solutions is finite - called mult. of  $\lambda_0$ .

(3) Assume  $\lambda_1, \lambda_2, \dots$  are roots of  $D(\lambda) = 0$  (listed acc to mult.)

$\phi_1(x), \phi_2(x), \dots$  ↙ orthonormal family of solns of  
 $\left( \int_a^b \phi_k(x) \phi_l(x) dx = \delta_{kl} \right)$  ↘  $\lambda_j \int K(x,y) \phi_j(y) dy = \phi_j(x)$

and assuming  $\sum_{k=1}^{\infty} \frac{\phi_k(x) \phi_k(y)}{\lambda_k}$  converges, we have

$$K(x,y) = \sum_{k=1}^{\infty} \frac{\phi_k(x) \phi_k(y)}{\lambda_k} \text{ and corresponding}$$

Fredholm det./resolvent are given by  $\prod_{k=1}^{\infty} \left(1 - \frac{\lambda}{\lambda_k}\right)$  (if convergent)

In this case, a series solution to Fredholm's equation is given as follows. Given  $f(x)$ , let

$$b_n = \int_a^b \phi_n(x) f(x) dx .$$

Then  $\phi(x) = \sum_{n=1}^{\infty} \frac{b_n \lambda_n}{\lambda_n - \lambda} \phi_n(x)$  solves

$$\phi(x) - \lambda \int_a^b K(x,y) \phi(y) dy = f(x)$$