

## Lecture 26

§0. Recall - we are studying integral equations (singular) and their relation with factorization problems. The crucial result in this setting is Plemelj formulae - stated below

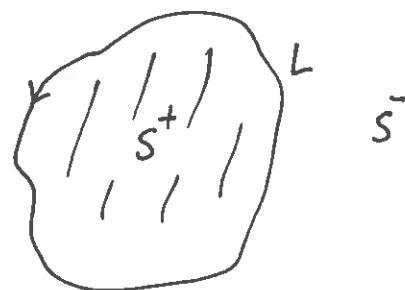
Let  $L$  be a contour in the complex plane,  $S^\pm$  connected components of (positively oriented)

$\mathbb{C} \setminus L$  ( $S^+$  is the bounded one - to the left of  $L$ )

Given  $\varphi : L \rightarrow \mathbb{C}$  satisfying

Hölder-type condition - we define

$$\Phi(z) := \frac{1}{2\pi i} \int_L \frac{\varphi(t)}{t-z} dt$$



Then  $\lim_{\substack{z \rightarrow t_0 \\ z \in S^\pm}} \Phi(z) = \pm \frac{1}{2} \varphi(t_0) + \frac{1}{2\pi i} \int_L \frac{\varphi(t)}{t-t_0} dt \quad \forall t_0 \in L$

↑  
Cauchy's principal value.

Thus, if  $\bar{\Phi}^\pm(t_0) := \lim_{\substack{z \rightarrow t_0 \\ z \in S^\pm}} \Phi(z)$ , then

$$\bar{\Phi}^+(t_0) - \bar{\Phi}^-(t_0) = \varphi(t_0)$$

$$\bar{\Phi}^+(t_0) + \bar{\Phi}^-(t_0) = \frac{1}{\pi i} \int_L \frac{\varphi(t)}{t-t_0} dt$$

(Plemelj formulae)

§1. Relation between singular integral equations and Riemann-Hilbert factorization problem.

$$(a) \quad a(t_0)\varphi(t_0) + \frac{b(t_0)}{\pi i} \int\limits_L \frac{\varphi(t)}{t-t_0} dt = f(t_0)$$

where  $a, b, f : L \rightarrow \mathbb{C}$  are given, and we have to solve for  $\varphi$ , is an example of a typical (dominant) singular integral equation.

Let  $\Phi(z) := \frac{1}{2\pi i} \int\limits_L \frac{\varphi(t)}{t-z} dt$  as before. By Plemelj formulae,

the singular integral equation given above, takes the following form

$$(a(t_0) + b(t_0)) \Phi^+(t_0) - (a(t_0) - b(t_0)) \Phi^-(t_0) = f(t_0)$$

an inhomogeneous R-H factorization problem.

(b) Conversely consider the (homogeneous) R-H factorization problem

$$\Phi^+(t_0) = G(t_0) \Phi^-(t_0)$$

$$\Phi^-(z) = \gamma(z) + O(z^{-1}) \text{ near } \infty$$

( $G$  and  $\gamma$  are given, and we want to find  $\Phi^\pm : S^\pm \rightarrow \mathbb{C}$ , which extend to cts. functions  $\pm$  on  $\overline{S^\pm}$ , whose values at  $t_0 \in L$  are related by  $G$ , as above).

By Cauchy's theorem, we obtain the following equations for  $\Phi^\pm$  - if they exist and extend to  $\partial S^\pm = L$  :

$$\frac{1}{2\pi i} \int_L \frac{\Phi^+(t) dt}{t-z} = \begin{cases} \bar{\Phi}^+(z) & \text{if } z \in S^+ \\ 0 & \text{if } z \in S^- \end{cases}$$

$$\frac{1}{2\pi i} \int_L \frac{\bar{\Phi}^-(t) dt}{t-z} - \gamma(z) = \begin{cases} -\bar{\Phi}^-(z) & \text{if } z \in S^- \\ 0 & \text{if } z \in S^+ \end{cases}$$

Combine with Plemelj formulae - we get

$$-\frac{1}{2} \bar{\Phi}^+(t_0) + \frac{1}{2\pi i} \int_L \frac{\bar{\Phi}^+(t) dt}{t-t_0} = 0$$

$$\frac{1}{2} \bar{\Phi}^-(t_0) + \frac{1}{2\pi i} \int_L \frac{\bar{\Phi}^-(t) dt}{t-t_0} - \gamma(t_0) = 0$$

Hence,  $\bar{\Phi}^+(t_0) = G(t_0) \bar{\Phi}^-(t_0)$  becomes a regular integral eq<sup>n</sup>:

$$\bar{\Phi}^-(t_0) - \frac{1}{2\pi i} \int_L \frac{G(t_0)^{-1} G(t) - I}{t-t_0} \bar{\Phi}^-(t) dt = \gamma(t_0)$$

Fredholm-type integral equation for  $\bar{\Phi}^-$ . Similarly for  $\bar{\Phi}^+$ :

$$\bar{\Phi}^+(t_0) - \frac{1}{2\pi i} \int_L \frac{I - G(t_0) G(t)^{-1}}{t-t_0} \bar{\Phi}^+(t) dt = G(t_0) \gamma(t_0).$$

§2. Fredholm theory. - Recall the main results from the previous lecture.

Given a cont. fn  $K(x,y)$ ,  $x,y \in [a,b]$ , we defined

$$D(\lambda) = 1 + \sum_{n=1}^{\infty} a_n \frac{\lambda^n}{n!} \quad \text{and} \quad D(x,y; \lambda) = \sum_{n=0}^{\infty} \frac{\lambda^{n+1}}{n!} v_n(x,y) \text{ where}$$

$$v_n(x,y) = (-1)^n \underbrace{\int_a^b \cdots \int_a^b}_{\text{n-fold}} \det \begin{bmatrix} K(x,y) & K(x,\xi_1) & \cdots & K(x,\xi_n) \\ K(\xi_1,y) & & & \\ \vdots & & & \\ K(\xi_n,y) & & & \end{bmatrix} d\xi_1 \cdots d\xi_n$$

$$\text{and } a_n = (-1)^n \underbrace{\int_a^b \cdots \int_a^b}_{\text{n-fold}} \det [K(\xi_i, \xi_j)]_{1 \leq i, j \leq n} d\xi_1, \dots d\xi_n$$

$$\text{Note: } a_n = - \int_a^b v_{n-1}(\xi, \xi) d\xi.$$

Theorem. - (i)  $D(\lambda)$  and  $D(x,y; \lambda)$  are entire fns. of  $\lambda$ .

(ii) For  $\lambda$  s.t.  $D(\lambda) \neq 0$ , there is a unique soln. to

$$\phi(x) - \lambda \int_a^b K(x,y) \phi(y) dy = f(x)$$

$$\text{given by } \phi(x) = f(x) + \frac{1}{D(\lambda)} \int_a^b D(x,y; \lambda) f(y) dy.$$

$$(iii) \int_a^b D(\xi, \xi; \lambda) d\xi = -\lambda \cdot D'(\lambda)$$

(iv) By col/row expansion of determinant, we get

$$\begin{aligned} D(x, y; \lambda) &= \lambda D(\lambda) K(x, y) + \lambda \int_a^b D(x, \xi; \lambda) K(\xi, y) d\xi \\ &= \lambda D(\lambda) K(x, y) + \lambda \int_a^b K(x, \xi) D(\xi, y; \lambda) d\xi \end{aligned}$$

(v) If  $D(\lambda_0) = 0$  to order  $\leq m$ , and  $D(x, y; \lambda)$  is expanded in Taylor series near  $\lambda_0$ :

$$D(x, y; \lambda) = g_\ell(x, y) (\lambda - \lambda_0)^\ell + \dots, \text{ then } m \geq \ell+1 \text{ and}$$

$$g_\ell(x, y) = \lambda_0 \int_a^b K(x, \xi) g_\ell(\xi, y) d\xi \text{ solves the hgr. eq^n}$$

$$\psi(x) = \lambda_0 \int_a^b K(x, \xi) \psi(\xi) d\xi, \quad (\text{i.e. } g_\ell(x, y_0) \text{ is a soln. of } \psi).$$

§3. Case of symmetric kernel (assume  $K$  is real-valued and  $K(x, y) = K(y, x)$ )

Theorem.  $D(\lambda) = 0$  has a (real) root  $\lambda_0$ .

(Schmidt)

Proof. Define  $K_n(x, y)$ ,  $n \geq 1$ , (Liouville's "iterated kernels")

$$\text{as: } K_1(x, y) = K(x, y)$$

$$K_{n+1}(x, y) = \int_a^b K(x, \xi) K_n(\xi, y) d\xi$$

Then, for  $|\lambda| < \frac{1}{M(b-a)}$  ( $M = \max \{|K(x, y)| : x, y \in [a, b]\}$ )

$$S(x) = f(x) + \lambda \int_a^b K_1(x, y) f(y) dy + \lambda^2 \int_a^b K_2(x, y) f(y) dy + \dots$$

converges uniformly and solves

$$S(x) - \lambda \int_a^b K(x,y) S(y) dy = f(x).$$

Now, let  $u_n = \int_a^b K_n(x,x) dx$ . Comparing the soln found in

Theorem §2 above with  $S(x) = f(x) + \lambda \int_a^b p(x,y;\lambda) f(y) dy$

where  $p(x,y;\lambda) = \sum_{n=0}^{\infty} K_{n+1}(x,y) \lambda^n$ , we have

$$D(x,y;\lambda) = \lambda D(\lambda) p(x,y;\lambda). \text{ Set } x=y=\xi \text{ and } \int_a^b,$$

use Thm §2 (iii) to get  $\boxed{\int_a^b p(\xi,\xi;\lambda) = -\frac{D'(\lambda)}{D(\lambda)}} \quad \text{i.e.,}$

$$\sum_{n=0}^{\infty} u_{n+1} \lambda^n = -\frac{D'(\lambda)}{D(\lambda)}.$$

Next, we show that  $u_2, u_4, \dots$  are all positive and

if  $v = \frac{u_4}{u_2}$ , then  $\frac{u_{2n+2}}{u_{2n}} \geq v^2$ . The first assertion is

clear since  $K$  is assumed to be symmetric and  $\mathbb{R}$ -valued

$$\begin{aligned} \text{e.g. } u_2 &= \int_a^b K_2(x,x) dx = \iint_{a,b} K(x,\xi) K(\xi,x) d\xi dx \\ &= \iint_{a,b} |K(x,\xi)|^2 dx d\xi > 0 \text{ and so on.} \end{aligned}$$

The second assertion follows from  $u_{2n+2} u_{2n-2} \geq u_{2n}^2$  which

can be proved by using

$$\int_a^b \int_a^b (\mu K_{n+1}(x, \xi) + K_{n-1}(x, \xi))^2 dx d\xi \geq 0 \quad (\forall \mu \in \mathbb{R})$$

$$\text{i.e. } \mu^2 u_{2n+2} + 2\mu u_{2n} + u_{2n-2} \geq 0.$$

In conclusion  $\frac{D'(\lambda)}{D(\lambda)} = - \sum_{n=0}^{\infty} u_{n+1} \lambda^n$  diverges for  $|\lambda|^2 \geq \frac{1}{n}$ . But

$D(\lambda)$  is an entire fn. - so is  $D'(\lambda)$  - hence  $D(\lambda) = 0$  for some

$\lambda_0$  within or on the circle  $|\lambda| = \sqrt{n^{-1/2}}$ .

□

§4. Continuing with the set up of §3, assume  $\lambda_0$  is s.t.  $D(\lambda_0) = 0$ .

Assume that there are  $n$  orthonormal solns. to

(say  $\psi_1, \dots, \psi_n$ )

$$\boxed{\begin{aligned} \lambda_0 \int K(x, y) \psi(y) dy \\ = \psi(x) \end{aligned}}$$

$$\text{We will show that } n \leq \lambda_0^2 \int_a^b \int_a^b K(x, y)^2 dx dy$$

hence there are only finitely many solns. to the hys. eq<sup>n</sup>

Proof. Let  $\Lambda(x, y) := \sum_{m=1}^n \left( \underbrace{\left( \int_a^b K(x, \xi) \psi_m(\xi) d\xi \right)}_{\frac{1}{\lambda_0} \psi_m(x) \text{ from the integral eq}^n} \right) \psi_m(y)$

$\frac{1}{\lambda_0} \psi_m(x)$  from the integral eq<sup>n</sup>.

Then

$$\int_a^b \Lambda(x, y)^2 dy = \sum_{m=1}^n \left( \int_a^b K(x, \xi) \psi_m(\xi) d\xi \right)^2$$

$$= \int_a^b K(x, y) \Lambda(x, y) dy$$

$$\Rightarrow \int_a^b \Lambda(x, y)^2 dy = * \int_a^b K(x, y)^2 - (K(x, y) - \Lambda(x, y))^2 dy$$

i.e..  $\int_a^b \left( \sum_{m=1}^n \frac{\psi_m(x) \psi_m(y)}{\lambda_0} \right)^2 dy \leq \int_a^b K(x, y)^2 dy$ . This means (by orthonormality)

$$\lambda_0^{-2} \sum_{m=1}^n \psi_m(x)^2 \leq \int_a^b K(x, y)^2 dy \text{ and the claim follows}$$

upon integrating w.r.t.  $x$ .

□