

# LECTURE 25

## Vector Field as a Flow

- I. Theorems: Where do they come from?
- II. Disassociated tangent spaces
- III. Vector field as a flow field
- IV. Flow field as a transformation
- V. Flow field as a transformation group

Read P125-127 in Chap. 5 of Singer and Thorpe's  
"Lecture Notes on Elementary Topology and Geometry"

Also read Sections 9.2-9.4 in MTW

I. Theorems: Where do they come from? or:

Is Deduction Possible without Induction?

Many admirers of mathematics, be they philosophers, engineers, physicists, including mathematicians themselves, hold mathematics in high regard as a deductive science. However, they are missing the point. Before one has deduction one must have induction.

Before one deduces that Socrates is a man, one must have induced (ultimately starting from observation) that all men are mortal.

Mathematical theorem are a good example. Before one proves a theorem (like the one on page 25.4), i.e. before one before one deduces certain consequences from what is known, given, or assumed, one must assemble, identify precisely, and state explicitly the starting point from which the subsequent consequences follow.

This assembling, identification, and articulation process - the inductive process - is much more difficult than a deduction. This is because induction is based on all of one's knowledge, a person's whole cognitive landscape.

Thus, because of limitations of time, resources, or even ignorance, perseverance and intelligence, instructors quite often skimp on the motivation, i.e. the necessary

inductive process which gives rise to the generalization expressed by the theorem, theory, principle, etc. This is done with the hope and understanding that the student's curiosity will pick up where the instructor did (or had to) slack off.

## II. Disassociated Tangent Spaces

Every point of a differentiable manifold  $M$  accommodates a tangent space which consists of vectors, i. e.

- a) of elements that correspond to  $n$ -tuples in  $\mathbb{R}^n$ , or
- b) of derivatives, or
- c) of tangents to curves in the manifold.

Thus, at each point of a differentiable manifold  $M$  there is a vector space space, the tangent space  $T_p(M)$  at  $p$ .

However, there is as yet no relation between all the vectors in  $T_p(M)$  and all those in  $T_{p'}(M)$  at a different points -

$P$  and  $P'$

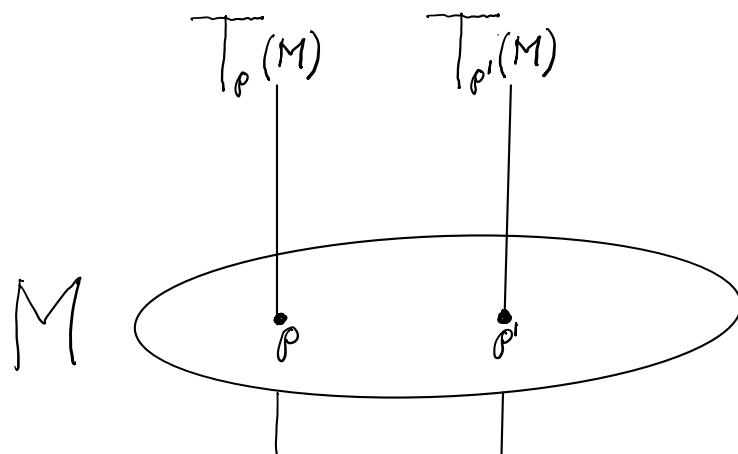


Figure 25.1: Tangent spaces  $T_p(M)$  and  $T_{p'}(M)$  are disjoint and are as yet unrelated.

Even with a given vector field  $u$ , i.e. an assignment of a vector  $u_p$  to each  $T_p(M)$ , there is no correspondence between  $T_p(M)$  and  $T_{p_1}(M)$ .

However, a vector field on  $M$  is a structure which does establish a correspondence between points on  $M$ , but not between the tangent spaces of a manifold.

### III. Vector Field as a Flow

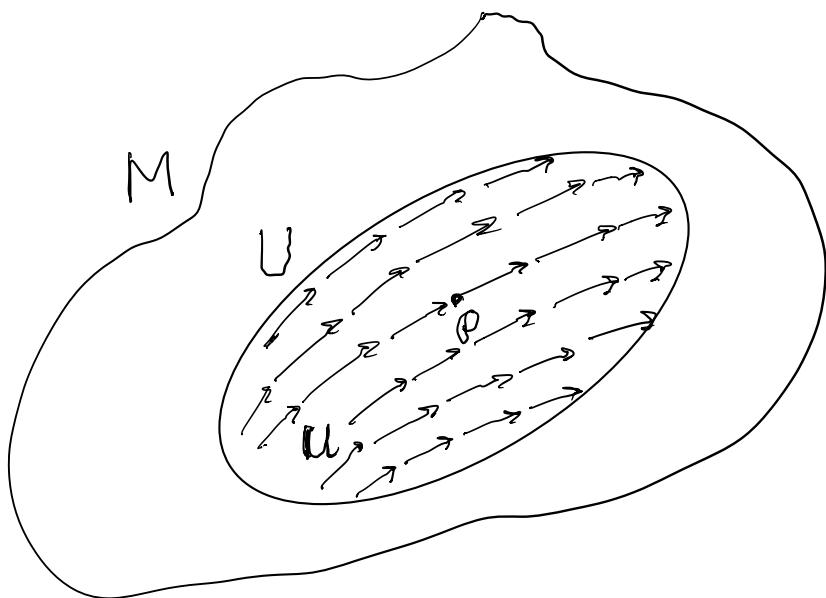


Figure 25.2: Vector field  $u$  in the neighborhood  $U$  of a point  $P$ .

A vector field induces a correspondence between points in a manifold. It is mathematized by the integral curves of the vector field viewed as a flow field. The existence of these curves is highlighted by the following

Definition ("Integral curve of a vector field")

Given: A smooth vector field (Figure 25.2)

$$u = u^i(x^1, \dots, x^n) \frac{\partial}{\partial x^i}$$

on M.

An integral curve of u is a smooth curve

$$c(\tau) : (a, b) \longrightarrow M$$

such that the vector  $\dot{c}(\tau)$  tangent to c at each point  $P \in c$  is the vector  $u(P)$  of the field u that has been assigned to  $T_p(M)$ :

$$\dot{c}(\tau) = u(c(\tau)) \quad a < \tau < b$$

or

$$\frac{d c^i}{d \tau} \frac{\partial}{\partial x^i} = u^i(c^1(\tau), \dots, c^n(\tau)) \frac{\partial}{\partial x^i}$$

or

$$\frac{d c^i}{d \tau} = u^i(c^1(\tau), \dots, c^n(\tau)) \quad i = 1, \dots, n$$

This is a system of n equations in n to-be-determined functions  $c^i(\tau)$ .

The existence and uniqueness of integral curves, i.e. of solutions to the system of 1<sup>st</sup> order o.d.e.'s is a feature of dynamical systems, including Hamiltonian systems, fluid dynamics, geodesic flows, in fact all systems of 1<sup>st</sup> order o.d.e.'s. Their features are captured by the following

Theorem ("Existence and Uniqueness")

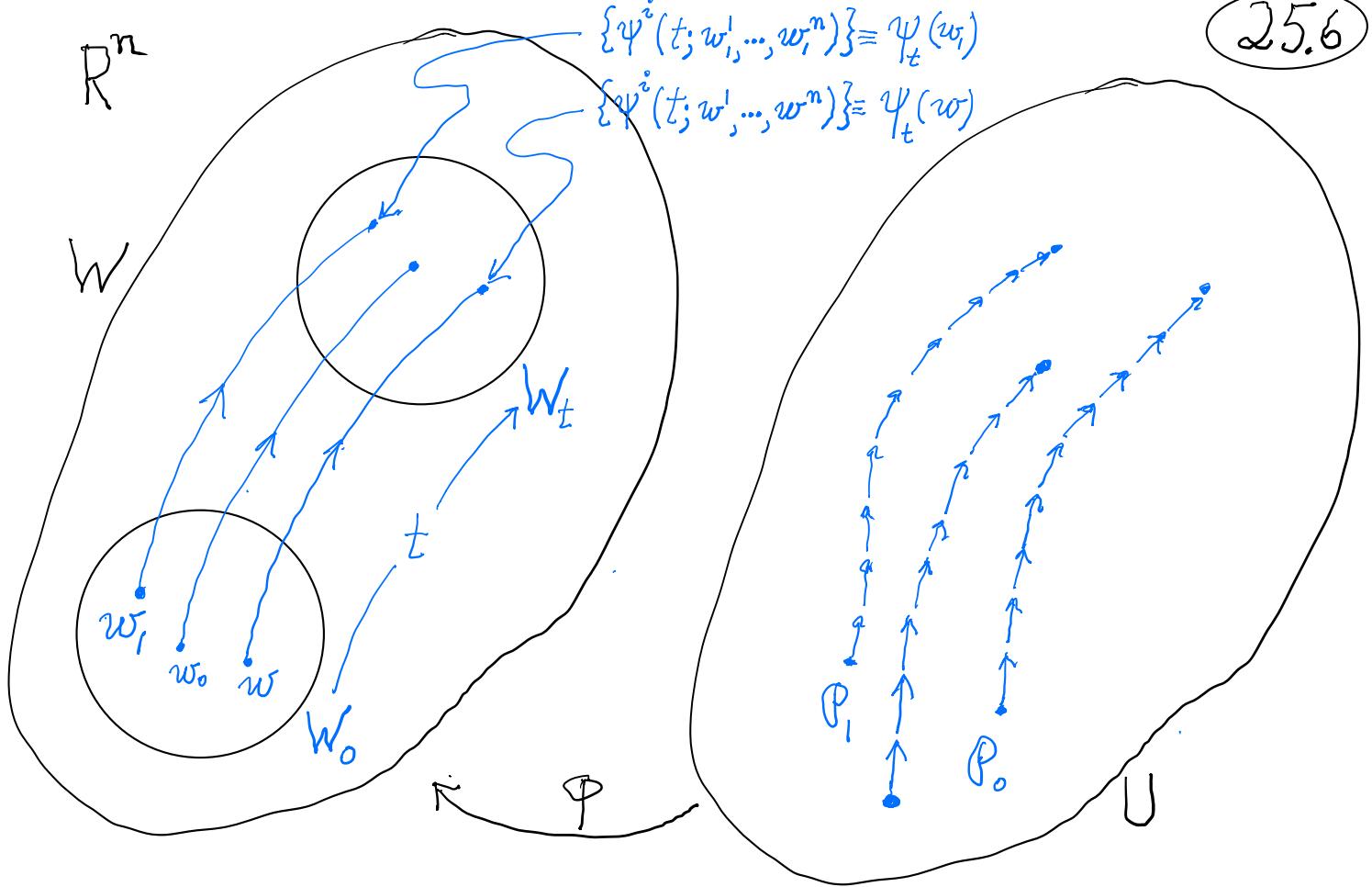


Figure 25.3: Integral curves of a vector field via their coordinate representation  $(\varphi, U)$ .

Given: Let  $W$  be an open set  $R^n$

Let  $w_0 = (w_0^1, \dots, w_0^n) \in R^n$

Let  $u^i(x^1, \dots, x^n) \in C^\infty(W, R)$ ,  $i=1, \dots, n$ , a set of functions smooth on  $W$

(Nota bene: They form the  $\varphi$ -representative of the vector field  $u = u^i \frac{\partial}{\partial x^i}$ )

Conclusion:

There exists

- (i) an open set  $W_0 \subset W$  about  $w_0$
- (ii) an open interval  $(-\epsilon, \epsilon) \subset R$
- (iii) a unique smooth  $n$ -component map

$$\psi = \{\psi^i\}_{i=1}^n : (-\epsilon, \epsilon) \times W_0 \longrightarrow W$$

$$(t; w^1, \dots, w^n) \mapsto \psi(t; w^1, \dots, w^n) = \{\psi^i(t; w^1, \dots, w^n)\}$$

such that  $\psi$  is a solution to the equation

$$\frac{d C^i(t)}{dt} = u^i(C^1(t), \dots, C^n(t)) \quad i = 1, \dots, n$$

subject to

$$C(0) = w, \text{ i.e. } C^i(0) = w^i, \quad i = 1, \dots, n$$

In other words, define the  $w = \{w^i\}$ -parametrized curve

$$\{C^i(t; w^1, \dots, w^n)\} = \{C_w^i(t)\}$$

$$\equiv C_w(t) : (-\epsilon, \epsilon) \longrightarrow W$$

by

$$\{C_w^i(t)\} = \{\psi^i(t; w)\} = \psi_t(w) \quad (\text{"flow field of } u\text{"})$$

with the property that

$$w = \{C_w^i(0)\} = \{\psi^i(0; w)\}. \quad (25.1)$$

Then for  $i = 1, \dots, n$  one has

$$\frac{d}{dt} [C_w^i(t)] = u^i(C_w^1(t), \dots, C_w^n(t))$$

#### IV. Flow Field as a Transformation

If  $C_w^i(t) = \psi^i(t; w)$  for  $i = 1, \dots, n$ , why does one use different symbols to refer to one and the same thing?

The reason is that the  $u$ -determined flow has two different aspects.

The symbol  $\{C_w^i\}$  is the reminder that the flow consists of individual curves that connect different points. By contrast, the notation

$$\{\psi^i(t; w)\} = \psi_t(w)$$

emphasizes the fact that  $\psi_t$  is a transformation that acts on the whole neighborhood  $W_0$  and transforms it into the new neighborhood  $W_t = \psi_t(W_0)$  at a different location without referring to the specific curves along which this happens.

## V. Flow Transformations form a Group under Composition

Recall that the flow of a vector field  $u$  is a  $w = \{w^1, \dots, w^n\}$  parametrized family of curves  $C_w(t) \equiv \psi_t(w)$  with tangents that are the preexisting vectors of the vector field.

The existence and uniqueness Theorem on pages 25.5-25.6 implies that every family member of the flow has the same additive property, namely  $\psi_{\tau} \circ \psi_{\tau'}(w) = \psi_{\tau+\tau'}(w)$ .

This property is brought to light by a notational shift in the flow function  $\psi$  as follows: Let

$$C_w(t) \equiv \psi_t(w) \equiv \{\psi^i(t; w^1, \dots, w^n)\}. \quad (25.2)$$

Then, as depicted in Figure 25.4,

$$\begin{aligned} \psi_t : W_0 &\longrightarrow W_t = \psi_t(W_0) \\ w &\rightsquigarrow \psi_t(w) \end{aligned} \quad (25.3)$$

is a  $t$ -parametrized group of one-to-one transformations that act on sets of points in the manifold.

- Let  $t = \tau' + \tau$  and apply the flow Theorem on pages 25.5-25.6 to the initial value data in  $W_0$ . The resulting transformation

depicted in Figure 25.4a is

$$\psi_{\tau'+\tau} : W_0 \rightarrow \psi_{\tau'+\tau}(W_0) = W_{\tau'+\tau} \quad (25.4)$$

$$w \mapsto w'' = \psi_{\tau'+\tau}(w) \in W_{\tau'+\tau} \quad (25.5)$$

25.9

This transformation exist and is unique for all  $-\epsilon < \tau', \tau < 0$ .

In particular, for  $\tau=0$  note that

$$W_{\tau'+0} = \psi_{\tau'+0}(W_0) \quad (25.6)$$

$$w' = \psi_{\tau'+0}(w). \quad (25.7)$$

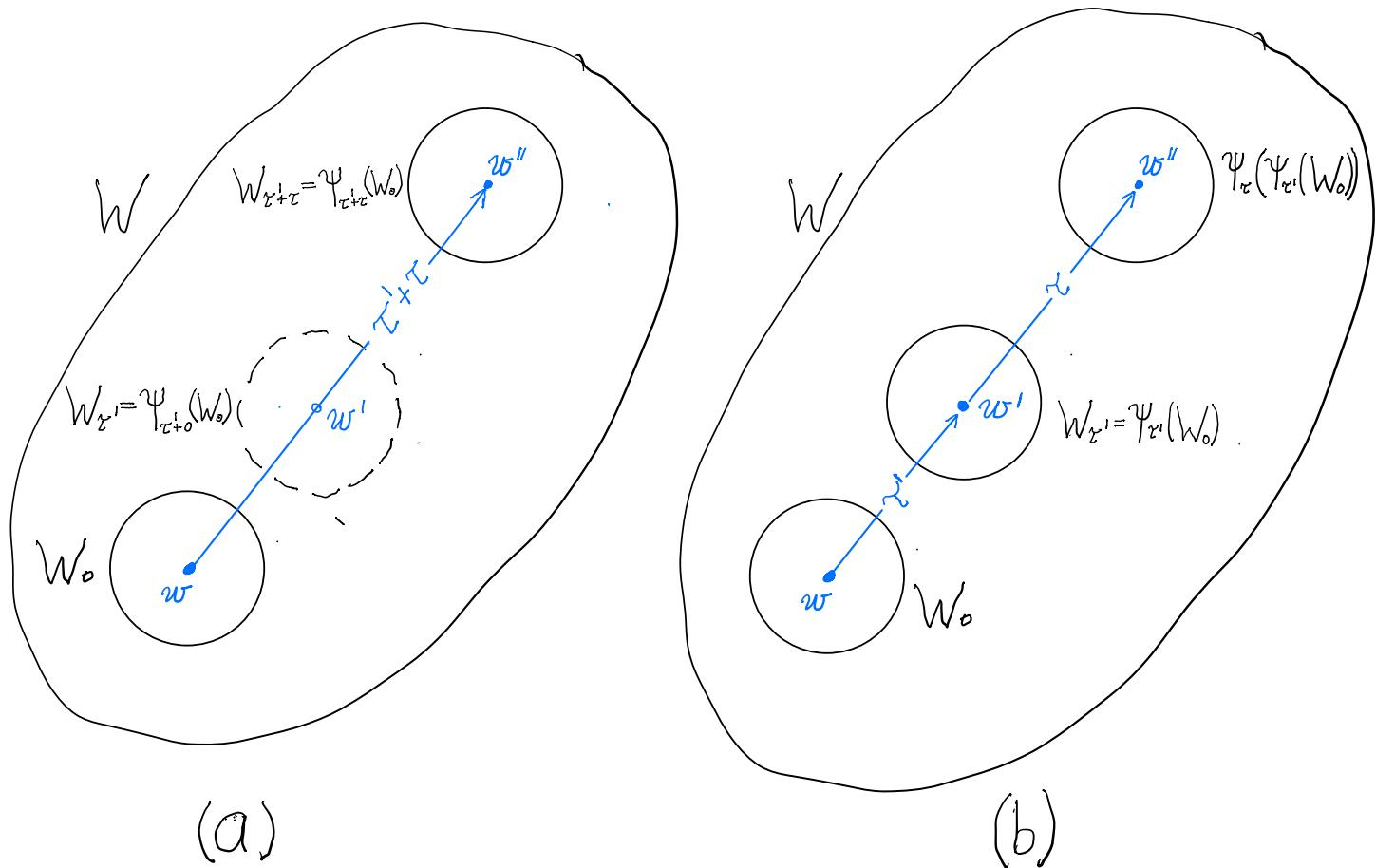


Figure 25.4: The  $u$ -induced transformation  $\psi_t$  applied to  $W_0$  in panel (a) results in  $\psi_{\tau+\tau}(W_0)$ . In panel (b) it results in  $\psi_{\tau'}\psi_{\tau}(W_0)$ . But both of them are also the transform images of one and the same  $\psi_{\tau'}(W_0)$ , the neighborhood centered around  $w'$ . Because of uniqueness the two are identical, i.e.  $\psi_{\tau+\tau}(W_0) = \psi_{\tau}\psi_{\tau'}(W_0)$ .

b) Let  $t = \tau'$  and apply the flow Theorem on pages 25.5-25.6 to the initial value data in  $W_0$ . The resulting transformation depicted in Figure 25.4b is

$$\psi_{\tau'} : W_0 \longrightarrow \psi_{\tau'}(W_0) = W_{\tau'} \quad (25.8)$$

$$w \rightsquigarrow w' = \psi_{\tau'}(w) \in W_{\tau'} \quad (25.9)$$

Using the points in the transformed neighborhood  $W_{\tau'}$ , Eq. (25.7), as initial value data, reapply to it the flow Theorem but with  $t = \tau$ . The resulting transformation depicted in Figure 25.4b is

$$\psi_{\tau} : \psi_{\tau'}(W_0) \longrightarrow \psi_{\tau}(\psi_{\tau'}(W_0)) = \psi_{\tau} \circ \psi_{\tau'}(W_0) \quad (25.10)$$

$$w' = \psi_{\tau'}(w) \rightsquigarrow w'' = \psi_{\tau}(\psi_{\tau'}(w)) = \psi_{\tau} \circ \psi_{\tau'}(w) \quad (25.11)$$

c) Compare Eqs. (25.10)-(25.11) with Eqs. (25.4)-(25.5) and infer that for all  $w$  in  $W_0$  one has

$$\psi_{\tau+\tau'}(w) = \psi_{\tau} \circ \psi_{\tau'}(w). \quad (25.12)$$

Indeed, the l.h.s. is according to Eq.(25.5) the image of  $\psi_{\tau'}(w)$ , Eq.(25.7), and the r.h.s. is according to Eq.(25.11) the image of the same  $\psi_{\tau'}(w)$ . Thus, because of the uniqueness of these images, the two sides are equal for all  $w \in W_0$ :

$$\boxed{\psi_{\tau+\tau'} = \psi_{\tau} \circ \psi_{\tau'}} \quad (25.13)$$

whose domain is

$$\psi_{\tau'}^{-1}(W_{\tau'}) = W_0 = \psi_{\tau+\tau'}^{-1}(W_{\tau''})$$

whenever  $-\epsilon < \tau, \tau', \tau + \tau' < \epsilon$ .

d) By letting  $\tau = -\tau'$  in the boxed Eq. (25.13) obtain

$$\psi_o(w) = \psi_{-\tau'} \circ \psi_{\tau'}(w). \quad (25.14)$$

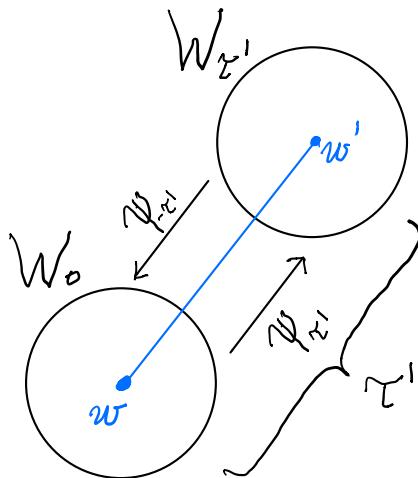


Figure 25.5: Flow transformation  $\psi_{-\tau'}$  composed with  $\psi_{\tau'}$  yields  $\psi_{-\tau'} \circ \psi_{\tau'} = \psi_o$ , the identity transformation:  $\psi_{-\tau'} = \psi_{\tau'}^{-1}$ .

But from the defining Eqs (25.2)-(25.3) and by using Eq.(25.1) in the flow Theorem on page 25.7, one has

$$\psi_o(w) = w \quad (25.15)$$

for all  $w \in W_0$ . Consequently, with Eq.(25.7) obtain

$$w = \psi_{-\tau'} \circ \psi_{\tau'}(w) \quad \forall w \in W_0$$

Thus,

$$\psi_{-\tau'} \circ \psi_{\tau'} = \text{identity}$$

and

$$\boxed{\psi_{-\tau'} = \psi_{\tau'}^{-1}} \quad (25.16)$$

# Conclusion

25.12

The two boxed Eqs. (25.13) and (25.16)

$$\boxed{\psi_{\tau+\tau'} = \psi_\tau \circ \psi_{\tau'}}$$

$$\boxed{\psi_{-\tau'} = \psi_{\tau'}^{-1}}$$

imply that the  $\psi$ 's form a local 1-parameter ( $\tau$ ) group of transformations.