

Lecture 30-Appendix

Holonomy: Rotation Generated by Curvature

Holonomy is rotation induced by parallel transport around closed curves 30A.1 with a common base point.

It holonomy is mathematized by a linear transformation which depends on the curvature integrated over the area bounded by a curve, say C . This linear transformation is given by

$$\begin{aligned} T_{\beta}^{\alpha}[C] &= \exp \oint_{\mathcal{S}} R_{\beta\mu\nu}^{\alpha}(\lambda, \eta) d\lambda d\eta \\ &= \exp \oint_{\mathcal{S}} R_{\beta\mu\nu}^{\alpha}(x^{\sigma}(\lambda, \eta)) d\lambda d\eta, \end{aligned} \quad (30A.1)$$

where \mathcal{S} is the area spanned by $x^{\sigma}(\lambda, \eta)$ and

$$\begin{aligned} u^{\rho} &= \frac{\partial x^{\rho}}{\partial \lambda} \\ v^{\sigma} &= \frac{\partial x^{\sigma}}{\partial \eta} \end{aligned}$$

are the component of the vectors $\frac{\partial}{\partial \lambda}$ and $\frac{\partial}{\partial \eta}$ tangent to \mathcal{S} .

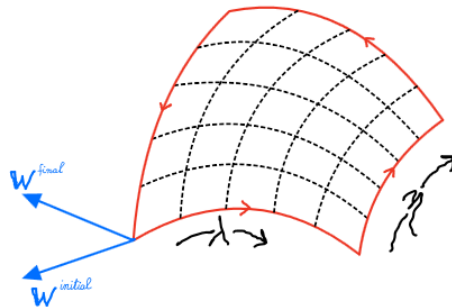


Figure 30A.1 When a vector w^{initial} is parallel transported around the boundary of the 2-d surface

$$\mathcal{S} = \{P(\lambda, \eta) : 0 \leq \lambda, \eta \leq 1\},$$

the specific nature of this parallel transport (as mathematized by its curvature tensor) transforms this vector into w^{final} .

30A.2

Where does the transformation Eq. (30A.1) come from, and how does one apply it?

1. Let $\mathcal{P}(\lambda, \eta): X^\alpha(\lambda, \eta)$ be a 2-d surface S spanned by λ and η .

2. Let $u = \frac{\partial}{\partial \lambda} \Big|_{\eta}$ and $v = \frac{\partial}{\partial \eta} \Big|_{\lambda}$

be the coordinate tangents to S .

3. Partition S into a union of small rectangles, each one being spanned by two small vectors

$$\Delta u_{ij} = \Delta \lambda \underbrace{\frac{\partial X^\alpha(\lambda, \eta)}{\partial \lambda}}_{u_{ij}^\alpha} \underbrace{\frac{\partial}{\partial X^\alpha}}_{e_\alpha} \equiv \Delta \lambda u_{ij}^\alpha e_\alpha$$

and

$$\Delta v_{ij} = \Delta \eta \underbrace{\frac{\partial X^\beta(\lambda, \eta)}{\partial \eta}}_{v_{ij}^\beta} \underbrace{\frac{\partial}{\partial X^\beta}}_{e_\beta} \equiv \Delta \eta v_{ij}^\beta e_\beta$$

Thus the (i, j) th rectangle is spanned by

$$(\Delta u_{ij}, \Delta v_{ij})$$

4. Let $\mathcal{P}(\lambda_i, \eta_j) \equiv \mathcal{P}_0(i, j)$ be a reference point on the boundary of the (i, j) th rectangle.

5. Let $W = w^x(\mathcal{P}) \frac{\partial}{\partial X^x} \equiv w^x(\mathcal{P}) e_x(\mathcal{P})$ be a vector field on S , and let

$$W(\mathcal{P}(\lambda_i, \eta_j)) = W_{ij}(\mathcal{P}_0) = e_x w_{ij}^x$$

be the vector value of W at $\mathcal{P}_0(i, j)$

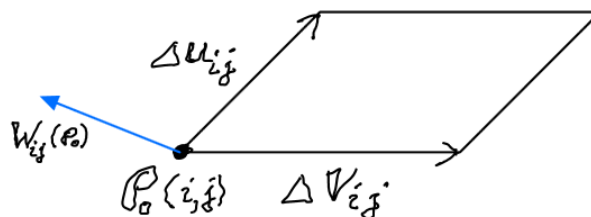


Figure 30A.2

30A.3

6. Let W_{ij}^{circum} be the result of parallel transporting $W_{ij}(P_0)$ around the perimeter of the $(i,j)^{th}$

rectangle:

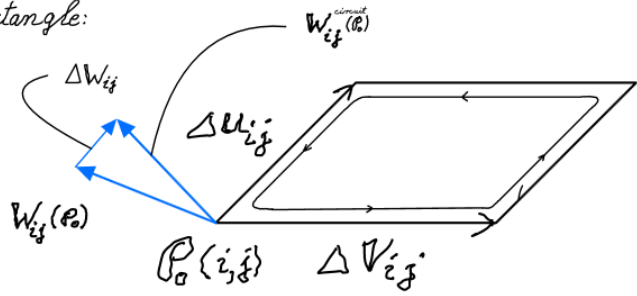


Figure 30A.3

Let ΔW_{ij} be the vectorial difference between the parallel transported vector and the one preexisting at $P_0(i,j)$:

$$\begin{aligned} \Delta W_{ij} &\equiv W_{ij}^{circum} - W_{ij}(P_0) = e_\alpha R^\alpha_\beta \Big|_{P_0(i,j)} (\Delta u_{ij}, \Delta v_{ij}) W_{ij}^\beta \\ &\equiv e_\alpha (d\omega^\alpha_\rho + \omega^\alpha_\gamma \wedge \omega^\gamma_\rho) \Big|_{P_0(i,j)} (\Delta u_{ij}, \Delta v_{ij}) W_{ij}^\beta \\ &= e_\alpha \underbrace{R^\alpha_{\beta\gamma\sigma} \Big|_{P_0(i,j)}}_{R^\alpha_{\beta uv}(i,j)} u_{ij}^\beta v_{ij}^\sigma \Delta\lambda \Delta\eta \end{aligned}$$

This difference is the change in the vector $W_{ij}(P_0)$ due to circumnavigational parallel transport around the boundary of S .

It follows that this circumnavigation implies

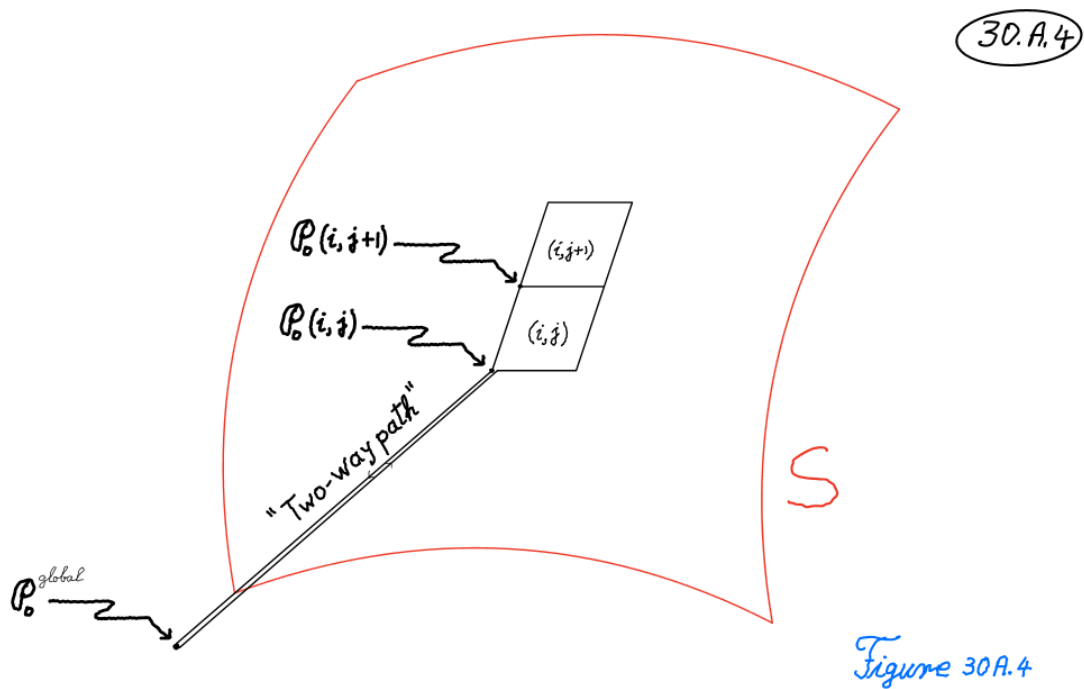
$$W_{ij}^{circum}(P_0) = e_\alpha(i,j) T^\alpha_\beta(i,j) W_{ij}^\beta$$

here

$$T^\alpha_\beta(i,j) = \delta^\alpha_\beta + R^\alpha_{\beta uv}(i,j) \Delta\lambda \Delta\eta$$

is a parallel transport induced linear transformation at $P_0(i,j)$. It differs infinitesimally from the identity transformation.

7. Let P_0^{glob} be a single global reference point on or off the boundary of S .



Let (i, j) and $(i, j+1)$ refer to two adjacent area elements of S .

Consider a two-way path

$$P_0^{global} \longleftrightarrow P_0(i, j)$$

which connects that single reference point with the point $P_0(i, j)$ at area element (i, j) .

8.a) Starting with some vector at P_0^{global} , say $w_{ij}^\beta(P_0^{global})$, parallel transporting it to $P_0(i, j)$, carrying it around the (i, j) th area element, and then back to P_0^{global} , one obtains

$$W_{ij}^{circuit}(P_0^{global}) = e_\alpha(P_0^{global}) T_\beta^\alpha(i, j) w_{ij}^\beta(P_0^{global}). \quad (30A.2)$$

$$(w_{i,j+1}^{circuit})^\alpha = T_{\beta}^{\alpha}(i,j+1) w_{i,j+1}^{\beta}$$

30A.6

$$= T_{\gamma}^{\alpha}(i,j+1) T_{\beta}^{\gamma}(i,j) w_{i,j}^{\beta}$$

9. Notice that this result is equivalent to starting with $w_{i,j}(\rho_0^{global})$ and circumnavigating both area elements (i,j) and $(i,j+1)$ at once. This is because the contributions to $w_{i,j+1}^{circuit}$ from adjacent sides cancel out.

10. In order to extend the process to all area elements, it is more economical to use matrix notation

$$[w_{i,j+1}^{circum}] = T(i,j+1) T(i,j) [w_{i,j}].$$

Applied to all area elements of S , one obtains

$$\begin{aligned} [w^{circum}] &= \prod_i \prod_j T(i,j) [w] \\ &= \prod_i \prod_j [I + R(i,j) \Delta\lambda \Delta\eta] [w] \end{aligned}$$

Here

$$[I + R(i,j) \Delta\lambda \Delta\eta]_{\beta}^{\alpha} = \delta_{\beta}^{\alpha} + R_{\beta\mu\nu}(i,j) \Delta\lambda \Delta\eta.$$

11. The product of the line integrals

$$\prod_i \prod_j [I + R(i,j) \Delta\lambda \Delta\eta] \quad (30A.3)$$

can be converted into an area integral over the two-dimensional area S . This is achieved by letting the partitioning of S into area elements become finer so that $\Delta u_{i,j}, \Delta v_{i,j}$ for the associated rectangles tend to zero.

To achieve such a conversion, take the limit of the logarithm of the matrix $(30A.7)$

product Eq. (30A.3):

$$\begin{aligned} & \lim_{\substack{\lambda_{i+1} \rightarrow \lambda_i \\ \eta_{j+1} \rightarrow \eta_j}} \log \prod_i \prod_j [I + R(i,j) \Delta \lambda \Delta \eta] & \Delta \lambda = \lambda_{i+1} - \lambda_i \\ & & \Delta \eta = \eta_{j+1} - \eta_j \\ & = \lim_{\substack{\lambda_{i+1} \rightarrow \lambda_i \\ \eta_{j+1} \rightarrow \eta_j}} \sum_i \sum_j \log [I + R(i,j) \Delta \lambda \Delta \eta] \\ & = \lim_{\substack{\lambda_{i+1} \rightarrow \lambda_i \\ \eta_{j+1} \rightarrow \eta_j}} \sum_i \sum_j R(i,j) \Delta \lambda \Delta \eta \end{aligned}$$

It follows that the in-the-limit infinite product is

$$\lim_{\substack{\lambda_{i+1} \rightarrow \lambda_i \\ \eta_{j+1} \rightarrow \eta_j}} \prod_i \prod_j [I + R(i,j) \Delta \lambda \Delta \eta] = \exp \int_S R(\lambda, \eta) d\lambda d\eta = [\exp \int_S R^*_{\beta\mu\nu}(\lambda, \eta) d\lambda d\eta]$$

and that the rotation experienced by a vector parallel transported around a loop that encloses the surface S is the linear transformation

$$W = e_\alpha W^\alpha \xrightarrow{\text{circum}} W = e_\alpha \exp \int_S R^*_{\beta\mu\nu}(\lambda, \eta) d\lambda d\eta W^\beta$$

CONCLUSION I

By starting with some vector w at some point event on a closed curve C , which is the boundary enclosing a surface S , one finds, upon parallel transporting w around C back to its starting point event, that

$$W \rightsquigarrow W^{\text{circum}}$$

is a linear transformation. It is mathematized by the surface 30A.8 integral of the curvature. The basis representation of this C -determined linear transformation is

$$T^\alpha_\beta [C] = \exp \oint_S R^\alpha_{\beta uv}(\lambda, \eta) d\lambda d\eta \quad (30A.4)$$

$$\equiv \exp \oint_S R^\alpha_{\beta \rho \sigma}(x^\alpha(\lambda, \eta)) u^\rho(\lambda, \eta) v^\sigma(\lambda, \eta) d\lambda d\eta$$

For metric compatible parallel transport

$$T^{\alpha\beta} [C] = \exp \oint_S R^{\alpha\beta}_{uv}(\lambda, \eta) d\lambda d\eta$$

is a rotation matrix. Thus, for every curve C there exists a rotation

transformation. This transformation is one which is orthonormal and the anti-symmetric matrix

$\oint_S R^{\alpha\beta}_{uv}(\lambda, \eta) d\lambda d\eta$ is its generator.

CONCLUSION I

It is easy to see that two curves can with the same starting point be strung together to form a third curve



The corresponding rotation transformation is

$$T[C_3] = T[C_2] \cdot T[C_1]$$

Thus, metric-compatible parallel transport establishes a mapping between closed curves and the group of Euclidean/Lorentz rotations.

CONCLUSION III

The agreement between rotation induced by (a) parallel transport around closed curves with a fixed base point

and (b) by the exponentiation of the corresponding area integrals of curvature form implies that the

(30A.9)

physical basis of the concept "curvature" is necessitated by the observation that things get rotated and these rotations form a group — the Holonomy Group at a base point and whose elements are generated by the curvature field permeating the area of the loops.

Application to a Two-sphere

As an application consider the holonomy of a meridian circle on a two-sphere of radius a , whose is

$$ds^2 = a^2 d\theta^2 + a^2 \sin^2 \theta d\varphi^2$$

and whose curvature relative to the coordinate bases

(30A.10)

$$\left\{ \frac{\partial}{\partial \theta}, \frac{\partial}{\partial \varphi} \right\} \text{ \& \} \{ d\theta, d\varphi \}$$

has components

$$R^{\theta}_{\varphi\theta\varphi} = \sin^2\theta$$

$$R^{\theta}_{\theta\theta\varphi} = 0$$

$$R^{\varphi}_{\theta\theta\varphi} = g^{\varphi\varphi} R_{\varphi\theta\theta\varphi} = -g^{\varphi\varphi} g_{\theta\theta} R^{\theta}_{\varphi\theta\varphi} = -1$$

$$R^{\varphi}_{\varphi\varphi\varphi} = 0.$$

Its components relative to the orthonormal basis

$$\left\{ \frac{1}{a} \frac{\partial}{\partial \theta}, \frac{1}{a \sin \theta} \frac{\partial}{\partial \varphi} \right\} \text{ \& \} \{ a d\theta, a \sin \theta d\varphi \}$$

$$\text{are } \hat{R}^{\theta}_{\varphi\theta\varphi} = (g_{\theta\theta} g^{\varphi\varphi} g^{\theta\theta} g^{\varphi\varphi})^{1/2} R^{\theta}_{\varphi\theta\varphi} = \frac{1}{a^2}$$

$$\hat{R}^{\theta}_{\theta\theta\varphi} = 0$$

$$\hat{R}^{\varphi}_{\theta\theta\varphi} = \hat{R}_{\varphi\theta\theta\varphi} = -\hat{R}_{\theta\varphi\theta\varphi} = -\hat{R}^{\theta}_{\varphi\theta\varphi} = -\frac{1}{a^2}$$

$$\hat{R}^{\varphi}_{\varphi\theta\varphi} = 0$$

In order to calculate the rotation, Eq. (30A.4) page 30A.8, first evaluate the integral

$$\oint_S \hat{R}^k_{\ell mn} \hat{\omega}^m \wedge \hat{\omega}^n.$$

Doing so for the domain $S = \{ \theta, \varphi: 0 \leq \theta \leq \theta_0, 0 \leq \varphi < 2\pi \}$, which is the northern polar cap bounded by the meridian $\theta = \theta_0$, one finds

$$\oint_S \hat{R}^k_{\ell mn} \hat{\omega}^m \wedge \hat{\omega}^n = \int_0^{2\pi} \int_0^{\theta_0} \underbrace{\hat{R}^k_{\ell\theta\varphi}}_{[-10] \frac{1}{a^2}} a d\theta a \sin \theta d\varphi$$

$$= \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix} 2\pi(1 - \cos\theta_0) \equiv \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix} \alpha \equiv K\alpha \rightarrow K = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$$

$$\alpha = 2\pi(1 - \cos\theta_0)$$

30A.11

Second, the rotation, Eq. (30A.4) on page 30A.8, is therefore

$$\exp K\alpha = \sum_{n=0}^{\infty} K^n \frac{\alpha^n}{n!}.$$

$$\text{Using } K^2 = \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} = -I$$

$$K^3 = -K$$

$$K^4 = I,$$

one obtains

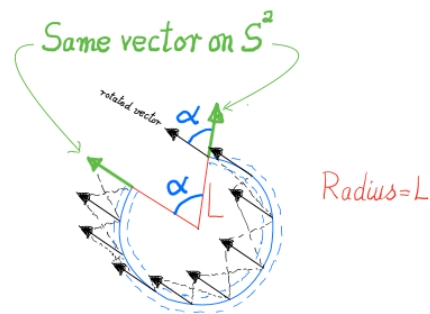
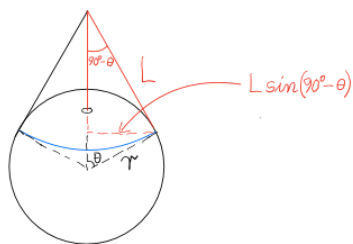
$$\begin{aligned} \exp K\alpha &= I \left(1 - \frac{\alpha^2}{2!} + \frac{\alpha^4}{4!} - \dots\right) + K \left(\alpha - \frac{\alpha^3}{3!} + \frac{\alpha^5}{5!} - \dots\right) \\ &= I \cos\alpha + K \sin\alpha \\ &= \begin{bmatrix} \cos\alpha & \sin\alpha \\ -\sin\alpha & \cos\alpha \end{bmatrix} \end{aligned}$$

Thus one has

$$\exp \oint \hat{R}^k_{\epsilon uv}(\lambda, \eta) d\lambda d\eta = \begin{bmatrix} \cos\alpha & \sin\alpha \\ -\sin\alpha & \cos\alpha \end{bmatrix}$$

where, as expected, $\alpha = 2\pi(1 - \cos\theta_0)$ is the angular rotation experienced by any vector parallel transported around the meridian $\theta = \theta_0$.

This experienced rotation can also be obtained from the following pictorial argument



30A.12

This angular rotation α due to parallel transport can also be established analytically by solving the equations for parallel transport

$$\nabla_v w = 0$$

which for $v = \dot{x}^\beta \frac{\partial}{\partial x^\beta}$ is equivalent to

$$\dot{w}^\alpha + \Gamma^\alpha_{\beta\gamma} \dot{x}^\beta w^\gamma = 0$$