

Lecture 33 & 34

Non-Euclidean Geometry in
the
Equatorial Plane of a Star

In MTW's chapter 23 read Section 23.8

I. Geometry of spacetime for a static star:

33.1

The spacetime geometry for any spherically symmetric system has the form

$$ds^2 = -e^{2\phi(r,t)} dt^2 + \frac{dr^2}{1 - \frac{2m(r,t)}{r}} + r^2(d\theta^2 + \sin^2\theta d\varphi^2)$$

For a system which is also static, there is no time dependence.

Its spatial geometry at any fixed time is therefore

$$ds^2 \Big|_{t=\text{fixed}} = \frac{dr^2}{1 - \frac{2m(r)}{r}} + r^2(d\theta^2 + \sin^2\theta d\varphi^2).$$

In the equatorial plane $\theta = \frac{\pi}{2}$ it is

$$ds^2 \Big|_{\substack{t=\text{fixed} \\ \theta = \frac{\pi}{2}}} \equiv ds^2 = \frac{dr^2}{1 - \frac{2m(r)}{r}} + r^2 d\varphi^2 \quad (\text{Non-Euclidean})$$

which is to be compared with

$$ds^2 = dr^2 + r^2 d\varphi^2 \quad (\text{Euclidean})$$

A. Imbedding Space

To obtain a geometrical picture of this non-Euclidean geometry, use the method of the imbedding diagram according to which one views the non-Euclidean plane as a surface of revolution in a 3-d fictitious imbedding space with a Euclidean geometry and spanned by its three coordinates z , r , and φ :

$$dL^2 = dz^2 + dr^2 + r^2 d\varphi^2 \quad (\text{metric for the imbedding space})$$

On the to-be-found surface of revolution

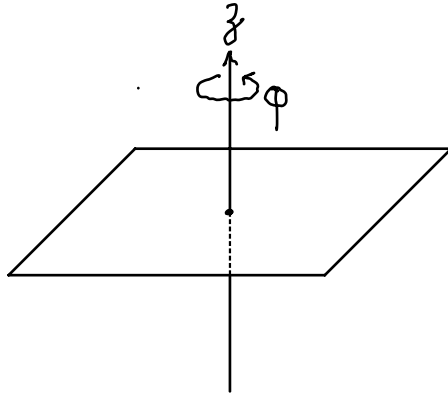


Figure 33.1 Fictitious 3-d imbedding space a non-Euclidean surface of revolution.

$$z = f(r)$$

The ambient Euclidean geometry induces the metric

$$dl^2 \Big|_{z=f(r)} = \left[\left(\frac{dz}{dr} \right)^2 + 1 \right] dr^2 + r^2 d\varphi^2$$

B. The Imbedding Function

Identify the metric on the to-be-found surface of revolution with the metric on the equatorial plane of the spherically symmetric spacetime. This results in the differential

equation

$$\left(\frac{dz}{dr} \right)^2 + 1 = \frac{1}{1 - \frac{2m(r)}{r}}$$

The solution to this differential equation

$$z(r) = \int_0^r \left[\frac{2m(r')}{r' - 2m(r')} \right]^{1/2} dr' + \text{const.} \quad (33.1)$$

yields a 2-d surface of revolution from a 3-d perspective.

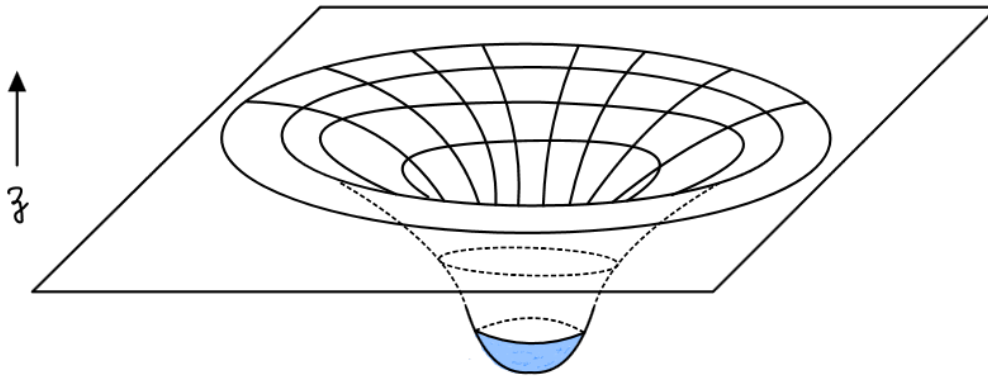


Figure 33.2 Imbedding diagram for the equatorial plane of a homogeneous star.

It allows one to visualize the inner 2-d spatial geometry on the equatorial plane or - because of spherical symmetry - any other rotated plane of the spherically symmetric space.

C. Example

Consider at some fixed time ($t = \text{const.}$) a star with mass density $\rho(r)$ in its interior and vacuum on the outside.

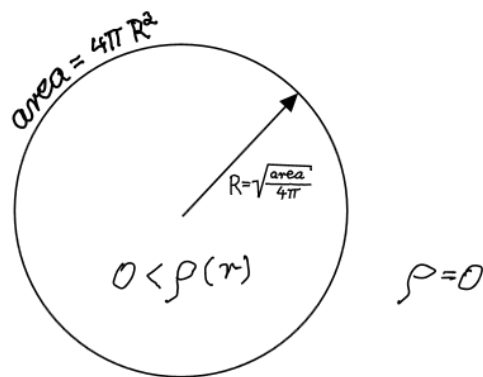


Figure 33.3 The radial parameter R for the concentric spheres of the star's interior and exterior is quantified in term of their proper area: $R = \sqrt{\text{area}/4\pi}$.

For such a configuration the mass function and its associated imbedding function are

$$m(r) = \begin{cases} \int_0^r 4\pi r'^2 \rho(r') dr' & \text{inside: } r < R \\ M & \text{outside: } r > R \end{cases}$$

and

$$z(r) = \begin{cases} \int_0^r \left[\frac{2m(r')}{r-2m(r')} \right]^{1/2} dr' + c & \text{inside: } r < R \quad (33.2) \\ [8M(r-2M)]^{1/2} + c & \text{outside: } r > R \quad (33.3) \end{cases}$$

Comment

Here the mass M and the mass density $\rho(r)$ are expressed in term of geometrical units:

$$M = \frac{G}{c^2} M^{\text{conventional}} = \left[\frac{G}{c^2} (\text{mass}) \right] = [\text{length}]$$

$$\rho = \frac{G}{c^2} \rho^{\text{conventional}} = \left[\frac{G}{c^2} \frac{(\text{mass})}{(\text{length})^3} \right] = \left[\frac{1}{(\text{length})^2} \right]$$

a) Thus outside the star one has

$$(z-c)^2 = 8M(r-2M)$$

which is a parabola of revolution.

b) Inside the star, near the center

$$m(r) = \frac{4\pi\rho_c}{3} r^3$$

The geometrized mass density has units $\frac{1}{(\text{length})^2}$.

Consequently, the density ρ_c implies a geometrically determined standard of length designated by \underline{a} :

$$\frac{8\pi}{3}\rho_c \equiv \frac{1}{a^2}.$$

With this the imbedding function in the central region inside the star has the form

$$\begin{aligned} z &= \int_0^r \sqrt{\frac{\left(\frac{r'}{a}\right)^2}{1-\left(\frac{r'}{a}\right)^2}} dr' = -a \sqrt{1-\left(\frac{r'}{a}\right)^2} \Big|_0^r \\ &= a - \sqrt{a^2 - r^2} \quad \text{for } r \ll a, \text{ near the center} \end{aligned}$$

Thus the imbedding function $z(r)$ is part of the circle of revolution:

$$(z-a)^2 + r^2 = a^2$$

c) At the star's boundary

$$\frac{dz}{dr} = \sqrt{\frac{2m(r)}{r-2m(r)}}$$

is continuous because $m(r)$ is continuous.

The geometry of a star is therefore characterized by a circle of revolution near its center, and a parabola of revolution outside its interior joined to its surface $r=R$

without any kinks. This is because $m(r)$ is continuous 33.6
there. Figure 33.2 depicts via its imbedding
diagram the equator plane and its non-Euclidean
geometry for a homogeneous star.